## Exercise 1. Composing Systems: The Tensor Product.

You have learned from quantum mechanics that the composition of two systems described by states  $|\psi_A\rangle \in \mathscr{H}_A$  and  $|\psi_B\rangle \in \mathscr{H}_B$  is described by a state in the tensor product space  $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle \in \mathscr{H}_A \otimes \mathscr{H}_B$ . The tensor product space is defined by its basis elements: if  $\{|\phi_A^i\rangle\}$  and  $\{|\phi_B^j\rangle\}$  are bases of  $\mathscr{H}_A$  and  $\mathscr{H}_B$ , respectively, then

$$\mathscr{H}_A \otimes \mathscr{H}_B = \operatorname{span}\left\{ |\phi_A^i\rangle \otimes |\phi_B^j\rangle \right\} \ . \tag{1}$$

The tensor product satisfies the following basic properties:

 $(|\psi_A\rangle + |\psi'_A\rangle) \otimes |\psi_B\rangle = |\psi_A\rangle \otimes |\psi_B\rangle + |\psi'_A\rangle \otimes |\psi_B\rangle ; \qquad (2)$ 

$$(\alpha |\psi_A\rangle) \otimes |\psi_B\rangle = \alpha \cdot |\psi_A\rangle \otimes |\psi_B\rangle , \qquad (3)$$

and the same properties hold on the second term.

- (a) Consider two qubits,  $\mathscr{H}_A = \mathscr{H}_B = \mathbb{C}^2$ , with respective bases  $\{|0_A\rangle, |1_A\rangle\}$  and  $\{|0_B\rangle, |1_B\rangle\}$ . The tensor product space admits the basis  $\{|0_A\rangle \otimes |0_B\rangle, |0_A\rangle \otimes |1_B\rangle, |1_A\rangle \otimes |0_B\rangle, |1_A\rangle \otimes |1_B\rangle\}$ . Write the state of each of the following systems in this basis.
  - (i) System A in state  $|0_A\rangle$  and system B in state  $|1_B\rangle$ .
  - (ii) System A in state  $\frac{1}{\sqrt{2}}(|0_A\rangle + |1_A\rangle)$  and system B in state  $|1_B\rangle$ .
  - (iii) System A in state  $\frac{1}{\sqrt{2}}(|0_A\rangle + |1_A\rangle)$  and system B in state  $\frac{1}{\sqrt{2}}(|0_B\rangle + |1_B\rangle)$ .

An important property of the tensor product space is that there are states in  $\mathscr{H}_A \otimes \mathscr{H}_B$  which cannot themselves be written as a tensor product of states from each space, i.e. they cannot be written in the form  $|\psi_A\rangle \otimes |\psi_B\rangle$ .

(b) Consider two qubits,  $\mathscr{H}_A = \mathscr{H}_B = \mathbb{C}^2$ , with respective bases  $\{|0_A\rangle, |1_A\rangle\}$  and  $\{|0_B\rangle, |1_B\rangle\}$ . Consider the state

$$|\Phi^{+}\rangle = \frac{1}{\sqrt{2}} \left(|0_{A}\rangle \otimes |0_{B}\rangle + |1_{A}\rangle \otimes |1_{B}\rangle\right) . \tag{4}$$

Show that this state vector cannot be written as a tensor product of two individual state vectors in  $\mathcal{H}_A$  and  $\mathcal{H}_B$ .

(c) Show that the state given in (b) can be written as

$$|\Phi^{+}\rangle = \frac{1}{\sqrt{2}} \left(|+_{A}\rangle \otimes |+_{B}\rangle + |-_{A}\rangle \otimes |-_{B}\rangle\right) , \qquad (5)$$

with  $|\pm_A\rangle = \frac{1}{\sqrt{2}} \left( |0_A\rangle + |1_A\rangle \right)$ 

(d) (Extra question, with an introduction to tomography.) We have shown that no state vector can appropriately describe the system A from point (b). However, it can be described by a density operator. Determine the density operator  $\rho$  for that system by considering explicitly the probabilities of the outcomes of the measurements in the bases  $\{|0_A\rangle, |1_A\rangle\}$ ,  $\{|+_A\rangle, |-_A\rangle\}$ , and  $\{|+i_A\rangle, |-i_A\rangle\}$  (where  $|\pm i_A\rangle = \frac{1}{\sqrt{2}}(|0_A\rangle \pm i|1_A\rangle)$ ).

*Hints.* By "measuring in a specific basis", it is meant to measure an observable that is diagonal in that basis. Recall also that the probability for measuring the outcome  $|\phi\rangle$  is given by  $\langle \phi | \rho | \phi \rangle$ .

In the density operator formalism, everything stays the same: the composition of two systems described by density operators  $\rho_A \in \mathcal{S}(\mathscr{H}_A)$  and  $\rho_B \in \mathcal{S}(\mathscr{H}_B)$  respectively, is described by a density operator  $\rho \in \mathcal{S}(\mathscr{H}_A \otimes \mathscr{H}_B) = \mathcal{S}(\mathscr{H}_A) \otimes \mathcal{S}(\mathscr{H}_B)$ . The important difference, however, is that  $\rho$  is not necessarily  $\rho_A \otimes \rho_B$ . Moreover, in contrast to state vectors, whatever the state of the joint system is, one can always write down the density operator of one part of the joint system, called *reduced state* or *marginal state*. The reduced state is obtained by *partial trace*.

- (e) Write out the density operators for the following systems using basis elements of  $\mathscr{H}_A \otimes \mathscr{H}_B$  given in point (a). (Use the matrix notation for convenience.)
  - (i) Two qubits in the state  $|\Phi^+\rangle$  defined in point (b).
  - (ii) Two qubits that are randomly prepared either jointly in state  $|0_A\rangle \otimes |0_B\rangle$  or in the joint state  $|1_A\rangle \otimes |1_B\rangle$ , with probability 1/2 each.
  - (iii) (Greenberger-Horne-Zeilinger state or cat state.) Three qubits A, B, and C in the state described by the vector

$$|\text{GHZ}\rangle = \frac{1}{\sqrt{2}} \left( |0_A\rangle \otimes |0_B\rangle \otimes |0_C\rangle + |1_A\rangle \otimes |1_B\rangle \otimes |1_C\rangle \right) \ . \tag{6}$$

(iv) The N-qubit version of the GHZ state,

$$|\text{GHZ}_N\rangle = \frac{1}{\sqrt{2}} \left(|0_1\rangle \otimes \cdots \otimes |0_N\rangle + |1_1\rangle \otimes \cdots \otimes |1_N\rangle\right) .$$
 (7)

(v) The maximally entangled state between two systems A and B, of N qubits each. Let  $\{|i_{A/B}\rangle\}_i$  be a basis for each system. The state is given by

$$|\Psi_N\rangle = \frac{1}{\sqrt{2^N}} \sum_i |i_A\rangle |i_B\rangle .$$
(8)

- (f) Calculate the following reduced states from point (e) and give their density operators.
  - (1.) The reduced state of system (i) on qubit A (respectively on qubit B).
  - (2.) The reduced state of system (ii) on qubit A (respectively on qubit B).
  - (3.) The reduced state of the GHZ state (iii) on the two first qubits, A and B.
  - (4.) The reduced state of the N-qubit GHZ state (iv) on all but the last qubit, i.e. just tracing out the N-th qubit.
  - (5.) The reduced state of the maximally entangled state  $|\Psi_N\rangle$  of point (v) on party A. Hint. Factorize the state vector cleverly.
  - (6.) The reduced state of the maximally entangled state  $|\Psi_N\rangle$  on the k first qubits of A and B (i.e. tracing out the N k last qubits of A and B).