

Heisenberg ferromagnet

model: nearest-neighbor ferromagnetic coupling

$$\mathcal{H} = -J \sum_{\langle i,j \rangle} \hat{\vec{S}}_i \cdot \hat{\vec{S}}_j = -J \sum_{\langle i,j \rangle} \left[\hat{S}_i^z \hat{S}_j^z + \frac{1}{2} \{ \hat{S}_i^+ \hat{S}_j^- + \hat{S}_i^- \hat{S}_j^+ \} \right]$$

$$s = \frac{\hbar}{2}$$

$$= -\frac{J}{V} \sum_{\vec{q}} \gamma_{\vec{q}} \hat{\vec{S}}_{\vec{q}} \cdot \hat{\vec{S}}_{-\vec{q}} = -\frac{J}{V} \sum_{\vec{q}} \gamma_{\vec{q}} \left[\hat{S}_{\vec{q}}^z \hat{S}_{-\vec{q}}^z + \frac{1}{2} \{ \hat{S}_{\vec{q}}^+ \hat{S}_{-\vec{q}}^- + \hat{S}_{\vec{q}}^- \hat{S}_{-\vec{q}}^+ \} \right]$$

$$\left. \begin{aligned} \hat{\vec{S}}_i &= \frac{1}{V} \sum_{\vec{q}} \hat{\vec{S}}_{\vec{q}} e^{i\vec{q} \cdot \vec{r}_i} \\ \gamma_{\vec{q}} &= 2 \sum_{\alpha=x,y,\dots} \cos(q_{\alpha}) \end{aligned} \right\}$$

conserved quantity

$$\hat{\vec{S}}_{tot} = \sum_i \hat{\vec{S}}_i = \hat{\vec{S}}_{\vec{q}=0}$$

mean field T_c

$$k_B T_c = J z \hbar^2$$

periodic boundary conditions

$$\vec{q} = \frac{2\pi}{L} (n_x, n_y, n_z)$$

$$[\hat{S}_i^z, \hat{S}_j^{\pm}] = \pm \hbar \delta_{ij} \hat{S}_i^{\pm}$$

$$[\hat{S}_i^+, \hat{S}_j^-] = 2\hbar \delta_{ij} \hat{S}_i^z$$

$$[\hat{S}_{\vec{q}}^z, \hat{S}_{\vec{q}'}^{\pm}] = \pm \hbar \hat{S}_{\vec{q}+\vec{q}'}^{\pm}$$

$$[\hat{S}_{\vec{q}}^+, \hat{S}_{\vec{q}'}^-] = 2\hbar \hat{S}_{\vec{q}+\vec{q}'}^z$$

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$$\chi(\vec{q}, t - t') = \frac{i}{\hbar} \Theta(t - t') \langle [\hat{S}_{\vec{q}}^-(t), \hat{S}_{-\vec{q}}^+(t')] \rangle .$$

equation of motion

$$i\hbar \frac{\partial}{\partial t} \chi(\vec{q}, t - t') = -\delta(t - t') \langle [\hat{S}_{\vec{q}}^-, \hat{S}_{-\vec{q}}^+] \rangle + \frac{i}{\hbar} \Theta(t - t') \langle [[\hat{S}_{\vec{q}}^-, \mathcal{H}](t), \hat{S}_{-\vec{q}}^+(t')] \rangle$$

$$= 2\delta(t - t') \hbar \langle \hat{S}_{\vec{q}=0}^z \rangle$$

$$- \frac{i}{\hbar} \Theta(t - t') \frac{J\hbar}{V} \sum_{\vec{q}'} (\gamma_{\vec{q}'} - \gamma_{\vec{q}+\vec{q}'}) \left\{ \left\langle \left[\hat{S}_{\vec{q}'}^z(t) \hat{S}_{\vec{q}-\vec{q}'}^-(t), \hat{S}_{-\vec{q}}^+(t') \right] \right\rangle \right. \\ \left. + \left\langle \left[\hat{S}_{\vec{q}+\vec{q}'}^-(t) \hat{S}_{-\vec{q}'}^z(t), \hat{S}_{-\vec{q}}^+(t') \right] \right\rangle \right\} .$$

Tyablikov decoupling

$$\frac{i}{\hbar} \Theta(t - t') \left\langle \left[\hat{S}_{\vec{q}'}^z(t) \hat{S}_{\vec{q}-\vec{q}'}^-(t), \hat{S}_{-\vec{q}}^+(t') \right] \right\rangle \rightarrow \begin{cases} \hat{S}_{\vec{q}'}^z(t) \hat{S}_{\vec{q}-\vec{q}'}^-(t) \rightarrow \langle \hat{S}_0^z \rangle \hat{S}_{\vec{q}}^-(t) \delta_{0,\vec{q}'} \\ \hat{S}_{\vec{q}+\vec{q}'}^-(t) \hat{S}_{-\vec{q}'}^z(t) \rightarrow \langle \hat{S}_0^z \rangle \hat{S}_{\vec{q}}^-(t) \delta_{0,\vec{q}'} \end{cases}$$

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decoupled equation of motion

$$i\hbar \frac{\partial}{\partial t} \chi(\vec{q}, t - t') = 2\delta(t - t')\hbar \langle \hat{S}_{\vec{q}=0}^z \rangle - \frac{J\hbar}{V} \langle \hat{S}_{\vec{q}=0}^z \rangle (\gamma_0 - \gamma_{\vec{q}}) \chi(\vec{q}, t - t')$$

Fourier transformation:

$$\int d\tilde{t} e^{i\omega\tilde{t} - \eta\tilde{t}} \left[i\frac{\partial}{\partial \tilde{t}} + \frac{J}{V} \langle \hat{S}_{\vec{q}=0}^z \rangle (\gamma_0 - \gamma_{\vec{q}}) \right] \chi(\vec{q}, \tilde{t}) = \int d\tilde{t} e^{i\omega\tilde{t} - \eta\tilde{t}} 2\delta(t - t') \langle \hat{S}_{\vec{q}=0}^z \rangle$$

$$\longrightarrow \left\{ \omega + i\eta + 2\frac{J}{V} \langle \hat{S}_0^z \rangle (\gamma_0 - \gamma_{\vec{q}}) \right\} \chi(\vec{q}, \omega) = 2\langle \hat{S}_0^z \rangle$$

$$\chi(\vec{q}, \omega) = \frac{2\langle \hat{S}_0^z \rangle}{\omega + 2\frac{J}{V} \langle \hat{S}_0^z \rangle (\gamma_0 - \gamma_{\vec{q}}) + i\eta}.$$

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spectrum of excitations $m \neq 0$

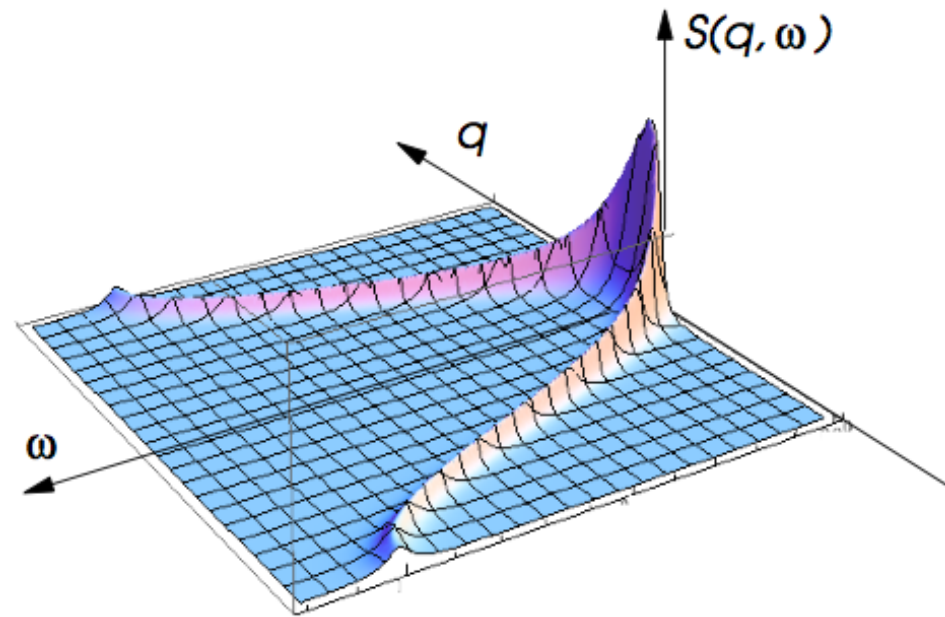
$$\chi''(\vec{q}, \omega) = 2\pi V m \delta[\omega - 2Jm(\gamma_0 - \gamma_{\vec{q}})]$$



$$\omega_{\vec{q}} = 2Jm(\gamma_0 - \gamma_{\vec{q}}) = 4Jm \sum_{\alpha} (1 - \cos q_{\alpha})$$

dynamical
structure factor

$$S(\vec{q}, \omega) = -\frac{1}{\pi} \frac{\chi''(\vec{q}, \omega)}{1 - e^{-\beta \hbar \omega}}$$



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instability condition

mean field $\langle \hat{S}_0^z \rangle = -mV$ $m(T \rightarrow T_{c-}) \rightarrow 0$ $k_B T_c^{\text{mf}} \approx 1.5 J \hbar^2$

$$S(\vec{q}, \omega) = \int dt e^{i\omega t} \frac{1}{\hbar} \langle \hat{S}_{\vec{q}}^-(t) \hat{S}_{-\vec{q}}^+(0) \rangle \quad \xrightarrow{\text{onsite equal-spin correlation function}}$$

$$\frac{1}{V^2} \sum_{\vec{q}} \int d\omega S(\vec{q}, \omega) = \frac{1}{\hbar} \langle \hat{S}_i^-(0) \hat{S}_i^+(0) \rangle = \frac{1}{\hbar} \left\{ \langle \hat{S}_i^2 \rangle - \langle \hat{S}_i^z \rangle^2 - \langle \hat{S}_i^z \rangle \hbar \right\} = \frac{\hbar}{2} + m$$

fluctuation-dissipation theorem

instability $m \rightarrow 0$

$$\frac{\hbar}{2} + m = -\frac{1}{\pi V^2} \sum_{\vec{q}} \int d\omega \frac{\chi''(\vec{q}, \omega)}{1 - e^{-\beta \hbar \omega}}$$

$$= \frac{1}{V} \sum_{\vec{q}} \frac{2m}{1 - e^{-\beta \hbar \omega_{\vec{q}}}}$$

determine $m(T)$

$$\frac{\hbar}{2} = \frac{k_B T_c}{J \hbar} \frac{1}{V} \sum_{\vec{q}} \frac{1}{\gamma_0 - \gamma_{\vec{q}}}$$

$d=3$

$$k_B T_c \approx 1.1 J \hbar^2$$

fluctuation renormiert

$d=1,2$

$$k_B T_c = 0$$

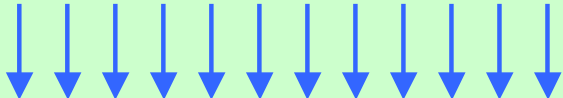
Hohenberg
Mermin-Wagner

spin wave in ferromagnet - magnon

Heisenberg ferromagnet

$$\mathcal{H} = -J \sum_{\langle i,j \rangle} \hat{\vec{S}}_i \cdot \hat{\vec{S}}_j = -J \sum_{\langle i,j \rangle} \left[\hat{S}_i^z \hat{S}_j^z + \frac{1}{2} \{ \hat{S}_i^+ \hat{S}_j^- + \hat{S}_i^- \hat{S}_j^+ \} \right]$$

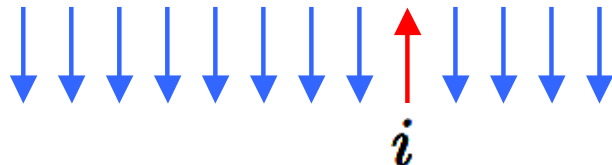
ground state:


 $|\psi^{(0)}\rangle = |\downarrow\downarrow \cdots \downarrow\downarrow\rangle$

$$\langle \hat{S}_i^z \rangle = -\frac{\hbar}{2}$$

$$\mathcal{H}|\psi^{(0)}\rangle = -\frac{J\hbar^2 N}{4}|\psi^{(0)}\rangle = E_0|\psi^{(0)}\rangle$$

excited state:

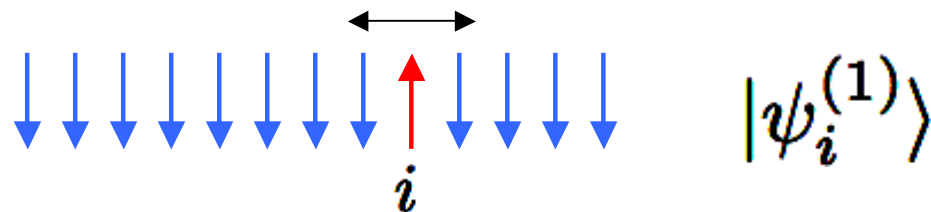

 $|\psi_i^{(1)}\rangle$

$$\mathcal{H}_1 = -J \sum_{\langle i,j \rangle} \hat{S}_i^z \hat{S}_j^z \quad \rightarrow \quad \mathcal{H}_1|\psi_i^{(1)}\rangle = E_1|\psi_i^{(1)}\rangle$$

$$E_1 = E_0 + \frac{Jz\hbar^2}{2}$$

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excited state: $\mathcal{H}_1 = -J \sum_{\langle i,j \rangle} \hat{S}_i^z \hat{S}_j^z \rightarrow \mathcal{H}_1 |\psi_i^{(1)}\rangle = E_1 |\psi_i^{(1)}\rangle$



$$E_1 = E_0 + \frac{Jz\hbar^2}{2}$$

$$\left. \begin{aligned} \mathcal{H}_2 &= -\frac{J}{2} \sum_{\langle i,j \rangle} \left\{ \hat{S}_i^+ \hat{S}_j^- + \hat{S}_i^- \hat{S}_j^+ \right\} \\ \hat{S}_i^+ |\downarrow\rangle &= \hbar |\uparrow\rangle \quad \hat{S}_i^- |\uparrow\rangle = \hbar |\downarrow\rangle \end{aligned} \right\} \rightarrow$$

"spin hopping"

$$\mathcal{H}_2 |\psi_i^{(1)}\rangle = -\frac{J\hbar^2}{2} \sum_{j \in \Lambda_i} |\psi_j^{(1)}\rangle$$

nearest neighbors of i

Fourier transformation $|\phi_{\vec{q}}^{(1)}\rangle = \frac{1}{V} \sum_i |\psi_i^{(1)}\rangle e^{-i\vec{q} \cdot \vec{r}_i}$ *"moving particle"*

$$\rightarrow \mathcal{H}_2 |\phi_{\vec{q}}^{(1)}\rangle = -\frac{J\hbar^2}{2} \sum_{j \in \Lambda_i} e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)} |\phi_{\vec{q}}^{(1)}\rangle = \frac{J\hbar^2}{2} \gamma_{\vec{q}} |\phi_{\vec{q}}^{(1)}\rangle$$

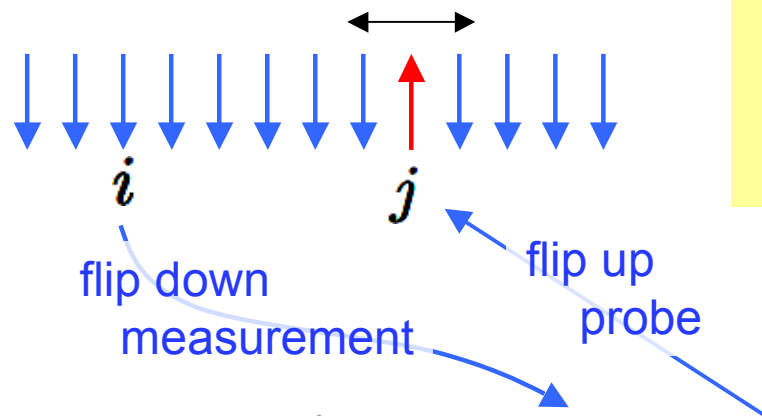
spin wave in ferromagnet - magnon

$$\left. \begin{aligned} \mathcal{H}_1 |\phi_{\vec{q}}^{(1)}\rangle &= \left(E_0 + \frac{Jz\hbar^2}{2} \right) |\phi_{\vec{q}}^{(1)}\rangle \\ \mathcal{H}_2 |\phi_{\vec{q}}^{(1)}\rangle &= \frac{J\hbar^2}{2} \gamma_{\vec{q}} |\phi_{\vec{q}}^{(1)}\rangle \end{aligned} \right\} \quad \mathcal{H} |\phi_{\vec{q}}^{(1)}\rangle = \left(E_0 + \frac{J\hbar^2}{2} (\gamma_0 - \gamma_{\vec{q}}) \right) |\phi_{\vec{q}}^{(1)}\rangle$$

$$|\phi_{\vec{q}}^{(1)}\rangle = \frac{1}{V} \sum_i |\psi_i^{(1)}\rangle e^{-i\vec{q} \cdot \vec{r}_i}$$

$$\gamma_{\vec{q}} = 2 \sum_{\alpha=x,y,z} \cos q_{\alpha} a$$

$$\Delta E_{\vec{q}} = J\hbar^2 \sum_{\alpha=x,y,z} (1 - \cos q_{\alpha} a)$$



$$\chi_{ij}(t) = \frac{i}{\hbar} \Theta(t) \langle [\hat{S}_i^-(t), \hat{S}_j^+(0)] \rangle$$

