Exercise 1. Oblique Corrections

In this exercise we want to investigate the so-called oblique corrections to electro-weak interactions. These corrections were introduced to analyse new physics phenomena for processes of $e^+e^- \rightarrow q\bar{q}$ as measured at LEP. In the Standard Model fermionic currents $J^{\mu}_Q(p^2)$ are coupled to the $SU_I(2) \times U_Y(1)$ gauge fields via the interaction Lagrangian

$$\mathcal{L} = gW^+_{\mu}J^{\mu}_{+} + gW^-_{\mu}J^{\mu}_{-} + gZ_{\mu}J^{\mu}_{Z} + eA_{\mu}J^{\mu}_{EM},$$

where

 $J^{\mu}_{Q} = i\bar{\psi}\gamma^{\mu}Q\psi.$

The different charges for the interactions are given by

$$Q_{\pm} = \frac{1}{\sqrt{2}} T^{\pm} P_L, \qquad Q_Z = \frac{1}{c_W} \left(T^3 - s_W^2 Q \right) P_L, \qquad Q_{EM} = eQ.$$

We want to consider one-loop corrections to resonant $q\bar{q}$ production in lepton collisions. Some leading corrections are given by fermion loops inserted into the gauge boson propagator. The gauge bosons A and B can be W^{\pm}, Z or γ .



We can decompose the resonant propagator corrections as

$$\Pi_{AB}^{\mu\nu}(p) = g^{\mu\nu}\Pi_{AB}(p^2) + \frac{p^{\mu}p^{\nu}}{p^2}\Delta_{AB}(p^2).$$

- (a) If we are interested in processes as described above why can we restrict ourselves to terms proportional to $g^{\mu\nu}$?
- (b) Calculated $\Pi_{AB}(p^2)$ for two left-handed $\Pi_{LL}(p^2)$ and one left- and one right-handed $\Pi_{LR}(p^2)$ current with two different fermion masses in the loop. Use dimensional regularisation in combination with anti-commuting γ^5 matrices. To compute the loop integrals introduce Feynman parameters and express your result in terms of the following integral definitions. Compute up to finite order in the dimensional regulator $d = 4 2\epsilon$.

$$b_0(p^2, m_1, m_2) = \int_0^1 dx \log \left(x m_1^2 + (1 - x) m_2^2 - p^2 x (1 - x) \right), \tag{1}$$

$$b_1(p^2, m_1, m_2) = \int_0^1 dx x \log \left(x m_1^2 + (1 - x) m_2^2 - p^2 x (1 - x) \right), \qquad (2)$$

$$b_2(p^2, m_1, m_2) = \int_0^1 dx x(1-x) \log \left(x m_1^2 + (1-x) m_2^2 - p^2 x(1-x) \right).$$
(3)

(c) Use your result from above to compute the correction to the photon propagator $e^2 \Pi_{QQ}(p^2)$. You should find that your corrections vanish as $p^2 \to 0$. Can you make a connection between the propagator correction of the photon and its renormalisation counter-terms? Why does the correction have to vanish as p^2 goes to zero?

In the lecture you introduced the Peskin-Takeuchi parameters S and T. These are related to the propagator corrections as

$$S \sim 4e^2 \left[\Pi'_{33}(0) - \Pi'_{3Q}(0) \right],$$
 (4)

$$T \sim \frac{e^2}{s_W^2 c_W^2 M_Z^2} \left[\Pi_{++}(0) - \Pi_{33}(0) \right].$$
 (5)

(d) Compute the leading corrections to the T parameter for the third family of standard model quarks. Consider the mass of the top quark to be larger than the bottom quark mass: $m_t >> m_b$. You should find

$$[\Pi_{++}(0) - \Pi_{33}(0)] = \frac{N_c}{16\pi^2} \frac{m_t^2}{4} + \mathcal{O}(\frac{m_b^2}{m_t^2}).$$