## Exercise 1. Bell-type Experiment.

Consider a 2-qubit Hilbert space $\mathscr{H}_{A B}=\mathscr{H}_{A} \otimes \mathscr{H}_{B}$ with basis $\{|00\rangle,|01\rangle,|10\rangle,|11\rangle\}$ in the Bell state

$$
\begin{equation*}
\left|\psi^{+}\right\rangle=\frac{1}{\sqrt{2}}\left(|0\rangle_{A}|0\rangle_{B}+|1\rangle_{A}|1\rangle_{B}\right) \tag{1}
\end{equation*}
$$

Two parties, Alice and Bob, get half of state $\left|\psi^{+}\right\rangle$so that Alice has qubit $A$ and Bob has qubit $B$. The POVM corresponding to a measurement can be written in function of the angle $\alpha$ that the measurement basis makes with the $\{|0\rangle,|1\rangle\}$ basis,

$$
M^{\alpha}=\left\{|\alpha\rangle\langle\alpha|,\left|\alpha^{\perp}\right\rangle\left\langle\alpha^{\perp}\right|\right\}, \quad|\alpha\rangle=\cos \frac{\alpha}{2}|0\rangle+\sin \frac{\alpha}{2}|1\rangle, \quad\left|\alpha^{\perp}\right\rangle=-\sin \frac{\alpha}{2}|0\rangle+\cos \frac{\alpha}{2}|1\rangle,
$$

where the $1 / 2$ factor comes from the Bloch sphere notation. We label the outcomes + for $|\alpha\rangle$ and - for $\left|\alpha^{\perp}\right\rangle$.

Suppose Alice performs such a measurement $\mathcal{M}^{\alpha}$ on her qubit.
(a) Find the description Bob would give to his partial state on $B$ after he knows that Alice performed the measurement $M_{A}^{\alpha}$ on $A$. What description would Alice give to $\rho_{B}$ given that she knows what measurement outcome she received?
(b) If Bob does the measurement $\mathcal{M}_{B}^{0}=\{|0\rangle\langle 0|,|1\rangle\langle 1|\}$ on $B$, what is the probability distribution for his outcomes, $\operatorname{Pr}_{B}$ ? How would Alice describe his probability distribution, $\operatorname{Pr}_{B \mid A}$ ?
(c) In part a) and b) Alice and Bob have different descriptions of the quantum state $\rho_{B}$ and probability distribution of measurement outcomes on that state. Explain how this subjective assignment of the scenarios at $B$ does not contradict with the actual measurement outcomes Bob will get after doing the measurement $\mathcal{M}_{B}^{0}$.

From now on look at the case where Alice and Bob can choose two different bases each:

(d) The joint probabilities $P_{X Y \mid a b}(x, y)$ of Alice and Bob obtaining outcomes $x$ and $y$ when they measure $A=a$ and $B=b$ are given in the following table.

| Alice |  | $\mathrm{A}=0$ |  | $\mathrm{~A}=2$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Bob |  | + | - | + | - |
| $\mathrm{B}=1$ | + | $\frac{1}{2}-\epsilon$ | $\epsilon$ | $\frac{1}{2}-\epsilon$ | $\epsilon$ |
|  | - | $\epsilon$ | $\frac{1}{2}-\epsilon$ | $\epsilon$ | $\frac{1}{2}-\epsilon$ |
| $\mathrm{B}=3$ | + | $\epsilon$ | $\frac{1}{2}-\epsilon$ | $\frac{1}{2}-\epsilon$ | $\epsilon$ |
|  | - | $\frac{1}{2}-\epsilon$ | $\epsilon$ | $\epsilon$ | $\frac{1}{2}-\epsilon$ |

$$
\text { with } \epsilon=\frac{1}{2} \sin ^{2}(\pi / 8) \approx 0.07
$$

Compute

$$
I_{N}\left(P_{X Y \mid A B}\right)=P(X=Y \mid A=0, B=3)+\sum_{|a-b|=1} P(X \neq Y \mid A=a, B=b)
$$

This quantity, similar to a Bell inequality, captures non-locality of quantum correlations: classically, it is at least equal to 1 , whereas for quantum correlations it can be smaller than 1.
(e) Correlations of the above form that exist within quantum theory cannot be created classically. However, they are not the most general distributions we could consider if we are only contained by the no-signalling principle: there are in fact other joint distributions that cannot be obtained by measurements on a quantum state, but that nonetheless would not allow for instantaneous information transmission over distance (signalling). To see this, look at the following joint probability distribution, a so-called PR box:

| Alice <br> Bob |  | $\mathrm{A}=0$ |  | $\mathrm{~A}=2$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| + | - | + | - |  |  |
| $\mathrm{B}=1$ | + | $\frac{1}{2}$ | 0 | $\frac{1}{2}$ | 0 |
|  | - | 0 | $\frac{1}{2}$ | 0 | $\frac{1}{2}$ |
| $\mathrm{~B}=3$ | + | 0 | $\frac{1}{2}$ | $\frac{1}{2}$ | 0 |
|  | - | $\frac{1}{2}$ | 0 | 0 | $\frac{1}{2}$ |

Show that the PR box
(i) is non-signalling: $P\left(X \mid a, b_{1}\right)=P\left(X \mid a, b_{2}\right), \forall a$;
(ii) is non-local: $P_{X Y \mid a b} \neq P_{X \mid a} P_{Y \mid b}$;
(iii) yields $I_{N}\left(P_{X Y \mid A B}\right)=0$.
(f) We shall now see how the above quantum correlation (coming from the Bell state) can be simulated using such a PR box combined with deterministic strategies. Imagine that Alice and Bob are given:

- with probability $1-p$ a PR-box;
- with probability $p / 4$, one of four deterministic boxes, that always outcome,+++- , -+ and -- respectively.

Find $p$ so that the final joint probability distribution equals the one of the Bell state given above.

Extra: Suppose that Alice and Bob are allowed to perform more measurements with closer angles (so that the overlap between two consecutive bases is larger). What happens to $p$ ?

See also: J. S. Bell, Phys. Vol. 1, No. 3, pp. 195-200, 1964.

