

**Exercise 11.1 Resource inequalities: teleportation and classical communication**

We saw a protocol, teleportation, to transmit one qubit using two bits of classical computation and one ebit,  $\frac{2}{\rightsquigarrow} \geq \frac{1}{\rightsquigarrow}$  (Section 6.1 of the script). Now suppose that Alice and Bob share unlimited entanglement: they can use up as many ebits as they want. Can Alice send  $n$  qubits to Bob using less than  $2n$  bits of classical communication? In other words, we want to know if the following is possible:

$$\frac{m}{\rightsquigarrow} \geq \frac{n}{\rightsquigarrow}, \quad m < 2n.$$

Prove that this is not the case. **Hint:** Use superdense coding.

**Exercise 11.2 A sufficient entanglement criterion**

In general it is very hard to determine if a state is entangled or not. In this exercise we will construct a simple entanglement criterion that correctly identifies all entangled states in low dimensions.

Recall that we say that a bipartite state  $\rho_{AB}$  is separable (not entangled) if

$$\rho = \sum_k p_k \sigma_k \otimes \tau_k, \quad \forall k : p_k \geq 0, \sigma_k \in \mathcal{S}_=(\mathcal{H}_A), \tau_k \in \mathcal{S}_=(\mathcal{H}_B), \quad \sum_k p_k = 1.$$

a) Let  $\Lambda_A : \text{End}(\mathcal{H}_A) \mapsto \text{End}(\mathcal{H}_A)$  be a positive map. Show that  $\Lambda_A \otimes \mathcal{I}_B$  maps separable states to positive operators.

This means that if we apply  $\Lambda_A \otimes \mathcal{I}_B$  to a bipartite state  $\rho_{AB}$  and obtain a non-positive operator, we know that  $\rho_{AB}$  is entangled. In other words, this is a sufficient criterion for entanglement.

b) Now we have to find a suitable map  $\Lambda_A$ . Show that the transpose,

$$\mathcal{T} \left( \sum_{ij} a_{ij} |i\rangle\langle j| \right) = \sum_{ij} a_{ji} |i\rangle\langle j|,$$

is a positive map from  $\text{End}(\mathcal{H}_A)$  to  $\text{End}(\mathcal{H}_A)$ , but is not completely positive.

c) Apply the partial transpose,  $\mathcal{T}_A \otimes \mathcal{I}_B$ , to the  $\varepsilon$ -noisy Bell state

$$\rho_{AB}^\varepsilon = (1 - \varepsilon) |\psi^-\rangle\langle\psi^-| + \varepsilon \frac{\mathbb{1}_4}{4}, \quad |\psi^-\rangle = \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle), \quad \varepsilon \in [0, 1].$$

For what values of  $\varepsilon$  can we be sure that  $\rho^\varepsilon$  is entangled?

Remark: Indeed, it can be shown that the PPT criterion (positive partial transpose) is necessary and sufficient for systems of dimension  $2 \times 2$  and  $2 \times 3$ .

**Exercise 11.3 Relative Entropy**

The quantum relative entropy is defined as  $D(\rho||\sigma) = \text{Tr}(\rho \log \rho - \rho \log \sigma)$ . For two classical probability distributions  $p$  and  $q$ , this definition simplifies to the expression for the Kullback-Leibler divergence  $\sum_j p_j \log \frac{p_j}{q_j}$ . Similar to the classical case, the relative entropy serves as a kind of “distance” between quantum states (although it is not technically a metric). Show that

- $H(A|B) = -D(\rho||\mathbb{1}_A \otimes \rho_B)$
- $D(\rho||\sigma) \geq 0$ , with equality if and only if  $\rho = \sigma$
- $D(\rho||\sigma) \leq \sum_k p_k D(\rho_k||\sigma)$ , where  $\rho = p_1 \rho_1 + p_2 \rho_2$
- For any CPTPM  $\mathcal{E}$ ,  $D(\rho||\sigma) \geq D(\mathcal{E}(\rho)||\mathcal{E}(\sigma))$
- $D(\rho||\sigma)$  is not a metric. Show this by proving that it is not symmetric.