## Exercise 5.1 Purification

A decomposition of a state  $\rho_A \in \mathcal{S}(\mathcal{H}_A)$  is a (non-unique) convex combination of pure states  $\rho_A^x = |a_x\rangle\langle a_x|$  such that  $\rho_A = \sum_x \lambda_x \rho_A^x$ .

- a) Show that  $|\Psi\rangle = \sum_x \sqrt{\lambda_x} |a_x\rangle_A \otimes |b_x\rangle_B$  is a purification of  $\rho_A$  for any orthonormal basis  $\{|b_x\rangle_B\}_x$  of  $\mathcal{H}_{\mathrm{B}}$ .
- b) Show that any two purifications are related by a *local* isometry on the purifying system.
- c) Mixed states can be decomposed in many different ways, i.e. with respect to many different bases. We will show that, from a purification of a mixed state  $\rho_A$ , we can generate any decomposition  $\{\rho_A^x\}_x$  such that  $\rho_A = \sum_x \lambda_x \rho_A^x$  by performing measurements on the purifying system. This is sometimes called *steering*.

For  $\rho_A$  as defined above, and any purification  $|\Phi\rangle$  of  $\rho_A$  on  $\mathcal{H}_A \otimes \mathcal{H}_B$ , find a measurement on  $\mathcal{H}_B$ , described by operators  $\{M_B^x\}_x$ , such that

$$\lambda_x = \operatorname{Tr}\left[|\Phi\rangle\langle\Phi|(\mathbb{1}_A \otimes M_B^x)\right] \quad \text{and} \quad \rho_A^x = \frac{\operatorname{Tr}_B\left[|\Phi\rangle\langle\Phi|(\mathbb{1}_A \otimes M_B^x)\right]}{\lambda_x}.$$
(1)

In this picture  $\lambda_x$  is the probability of measuring x and  $\rho_A^x$  is the state after the measurement.

**Decomposition example.** To see what we mean by the different compositions, look at the following example for the fully mixed state. It can be written as

$$\begin{split} \frac{\mathbb{1}_2}{2} &= \frac{|0\rangle\langle 0| + |1\rangle\langle 1|}{2}, \quad \text{or} \\ &= \frac{|+\rangle\langle +| + |-\rangle\langle -|}{2}, \quad \text{or (more interestingly)} \\ &= \sum_{i=1}^4 \frac{1}{4} |\theta_i, \phi_i\rangle\langle \theta_i, \phi_i|, \end{split}$$

where  $\{|\theta_i, \phi_i\rangle\}$  are the pure states:

$$(\theta_1,\phi_1) = (\psi,\frac{\pi}{4}), \quad (\theta_2,\phi_2) = (\pi-\psi,-\frac{\pi}{4}), \quad (\theta_3,\phi_3) = (\pi-\psi,\frac{3\pi}{4}), \quad (\theta_4,\phi_4) = (\psi,-\frac{3\pi}{4}),$$

with  $\psi := \arccos(\frac{1}{\sqrt{3}})$ . Note that as these states are pure, they lie on the surface of the sphere, and therefore can be parametrized by just two parameters  $(\theta, \phi) \in [0, \pi] \times [0, 2\pi)$ . You can also see that the vectors sit on the vertices of a regular tetrahedron.

## Exercise 5.2 Distinguishing two quantum states

Suppose you know the density operators of two quantum states  $\rho, \sigma \in \mathcal{H}_A$ . Then you are given one of the states at random—it may either be  $\rho$ , with probability p, or  $\sigma$ , with probability 1 - p. The challenge is to perform a single measurement on your state and then guess which state that is.

- a) What is your best strategy? In which basis do you think you should perform the measurement? Can you express that measurement using a projector P?
- b) What is the probability of guessing correctly,  $\Pr^p_{\checkmark}(\rho, \sigma)$ ? Compare that with the case where the states are evenly distributed,  $\Pr^{0.5}_{\checkmark}(\rho, \sigma) = \frac{1}{2}[1 + \delta(\rho, \sigma)]$ , where  $\delta(\rho, \sigma)$  is the trace distance between the two quantum states.

## Exercise 5.3 Distance bounds

Maximally entangled states  $|\Psi\rangle$  between two systems A and A' are used for many quantum tasks (for instance teleportation, which we will introduce later in the lecture). Unfortunately, sometimes we cannot be 100% sure we can create exactly  $|\Psi\rangle$ . Suppose, however, that we do know how to create a state in  $A \otimes B$  such that B has almost no information about A:  $\rho_{AB} \approx \mathbb{1}_A/|A| \otimes \rho_B$ . As we shall see, this can help us find an (approximately) maximally entangled state between A and A'.

a) Given a trace-preserving quantum operation  $\mathcal{E}$  and two states  $\rho$  and  $\sigma$ , show that

$$\delta\left(\mathcal{E}(\sigma),\mathcal{E}(\rho)\right) \leq \delta(\sigma,\rho)$$

How would you interpret this statement?

b) Show that any purification of the state  $\rho_{AB} = \frac{\mathbb{1}_A}{|\mathcal{H}_A|} \otimes \rho_B$  has the form

$$|\psi\rangle_{AA'BB'} = |\Psi\rangle_{AA'} \otimes |\psi\rangle_{BB'},$$

where  $|\Psi\rangle_{AA'} = |\mathcal{H}_A|^{-\frac{1}{2}} \sum_i |i\rangle_A |i\rangle_{A'}$  is a maximally entangled state, and  $|\psi\rangle_{BB'}$  is a purification of  $\rho_B$ .

- c) Show that  $1 F(\rho, \sigma) \le \delta(\rho, \sigma) \le \sqrt{1 F(\rho, \sigma)^2}$ .
- d) Consider a state  $\sigma_{AB}$  that is  $\varepsilon$ -distant from  $\rho_{AB}$  according to the trace distance, i.e.

$$\delta\left(\sigma_{AB}, \frac{\mathbb{1}_A}{|\mathcal{H}_A|} \otimes \rho_B\right) \leq \varepsilon.$$

Find an upper bound for

$$\delta\left(\tau_{AA'}, |\Psi\rangle_{AA'} \langle \Psi|_{AA'}\right)$$

where  $|\phi\rangle_{AA'BB'}$  is a purification of  $\sigma_{AB}$  and  $\tau_{AA'} = \text{Tr}_{BB'}(|\phi\rangle_{AA'BB'}\langle\phi|_{AA'BB'})$ .