

1. Linear sigma model with trivial vacuum

Consider a model of N real scalar fields Φ^i that couple to each other through a quartic interaction that is symmetric under $SO(N)$ rotations of the N fields. The Lagrangian of this model is

$$\mathcal{L} = -\frac{1}{2}\partial^\mu\Phi^i\partial_\mu\Phi^i - \frac{1}{2}m^2\Phi^i\Phi^i - \frac{1}{8}\lambda(\Phi^i\Phi^i)^2. \quad (1)$$

a) Derive the corresponding Hamiltonian and show that

$$V(\Phi) = \frac{1}{2}m^2\Phi^i\Phi^i + \frac{1}{8}\lambda(\Phi^i\Phi^i)^2 \quad (2)$$

is the potential term of the Hamiltonian.

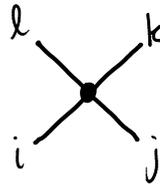
First consider the case $m^2 > 0$, convince yourself that for $\lambda = 0$ the Hamiltonian is just an N -fold copy of the Klein–Gordon Hamiltonian. For small λ we can calculate a perturbation series in λ .

b) Show that the Wick contraction of the Φ^i fields is

$$\underbrace{\Phi^i(x)\Phi^j(y)} = -i\delta^{ij}G_F(x-y), \quad (3)$$

where G_F is the Feynman propagator for a Klein-Gordon scalar field of mass m .

c) Show that there is one interaction vertex given by



$$= -i\lambda(\delta^{ij}\delta^{kl} + \delta^{il}\delta^{jk} + \delta^{ik}\delta^{jl}). \quad (4)$$

d) Let $N = 2$ and compute at leading order in λ the differential cross section $d\sigma/d\Omega$ for

$$\Phi^1\Phi^1 \rightarrow \Phi^2\Phi^2, \quad (5)$$

$$\Phi^1\Phi^2 \rightarrow \Phi^1\Phi^2, \quad (6)$$

$$\Phi^1\Phi^1 \rightarrow \Phi^1\Phi^1. \quad (7)$$

Note that the differential cross section for four particles in the centre of mass frame is given by

$$\frac{d\sigma}{d\Omega} = \frac{|M|^2}{64\pi^2s}, \quad (8)$$

where s is the centre of mass energy squared of the system and M is the matrix element describing the process.

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2. Linear sigma model with non-trivial vacuum

Next we consider the case where $m^2 =: -\mu^2 < 0$. Convince yourself that $V(\Phi)$ has a local maximum at $\Phi^i = 0$. As the potential is bounded from below, the minimum must be located at a non-vanishing value of Φ^i . Moreover, the theory is invariant under global $SO(N)$ rotations of the fields, and all points on the sphere with equal $|\Phi|$ must also be minima of $V(\Phi)$. The ground state of our field theory is therefore given by some non-zero constant field Φ^i . We choose Φ^i to point along the N -th direction or use a $SO(N)$ rotation to that end. We parametrise the quantum fields around the vacuum as

$$\Phi^i(x) = \phi^i(x), \quad i = 1, \dots, N-1, \quad (9)$$

$$\Phi^N(x) = v + \sigma(x). \quad (10)$$

- Determine the vacuum expectation value v , i.e. the field value in the minimum, in terms of μ and λ by minimising the potential $V(\Phi)$.
- Insert the ansatz for Φ in terms of ϕ and σ and the expression for v into the Lagrangian (or Hamiltonian) and show that the new Lagrangian (Hamiltonian) describes a theory of a massive field σ and $N-1$ massless fields ϕ^i .
- Convince yourself that the ϕ and σ fields interact through a new set of cubic and quartic vertices and determine the Feynman rules for all propagators and vertices.

$$\underbrace{\phi^i(x) \phi^j(y)} = \begin{array}{c} \times \\ \bullet \\ \text{---} \\ \bullet \\ \gamma \end{array} \quad \underbrace{\sigma(x) \sigma(y)} = \begin{array}{c} \times \\ \bullet \\ \text{---} \\ \bullet \\ \gamma \end{array} \quad (11)$$

$$\begin{array}{ccccc} \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ i \quad j \end{array} & \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ i \quad j \end{array} & \begin{array}{c} l \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ i \quad j \end{array} & \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ i \quad j \end{array} & \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ i \quad j \end{array} \end{array} \quad (12)$$

- Let $N = 2$ and compute at leading order in λ the differential cross section $d\sigma/d\Omega$ for

$$\phi^1 \phi^1 \rightarrow \phi^2 \phi^2, \quad (13)$$

$$\phi^1 \phi^2 \rightarrow \phi^1 \phi^2, \quad (14)$$

$$\phi^1 \phi^1 \rightarrow \phi^1 \phi^1. \quad (15)$$

Note that there are now four Feynman diagrams contributing to the amplitude at leading order

$$M = \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \\ \bullet \\ \text{---} \\ \bullet \\ \text{---} \\ \text{---} \end{array} \quad (16)$$