

**Exercise 7.1 Exact solution of the Ising chain**

In this exercise we will investigate the physics of one of the few *exactly solvable interacting* models, the one-dimensional Ising model (Ising chain). Consider a chain of  $N + 1$  Ising-spins with free ends and nearest neighbor coupling  $-J$  ( $J > 0$  for ferromagnetic coupling)

$$\mathcal{H}_{N+1} = -J \sum_{i=1}^N \sigma_i \sigma_{i+1}, \quad \sigma_i = \pm 1. \quad (1)$$

In this exercise we will be interested in the thermodynamic limit of this system, i.e. we assume  $N$  to be very large.

- Compute the partition function  $Z_{N+1}$  using a recursive procedure.
- Calculate the magnetization density  $m = \langle \sigma_j \rangle$  where the spin  $\sigma_j$  is far away from the ends. Which symmetries does the system exhibit? Interpret your result in terms of symmetry arguments.
- Show that the *spin correlation function*  $\Gamma_{ij} = \langle \sigma_i \sigma_j \rangle - \langle \sigma_i \rangle \langle \sigma_j \rangle$  decays exponentially with increasing distance  $|j - i|$  on the scale of the so-called *correlation length*  $\xi$ , i.e.  $\Gamma_{ij} \sim e^{-|j-i|/\xi}$ . Show that  $\xi = -[\log(\tanh \beta J)]^{-1}$  and interpret your result in the limit  $T \rightarrow 0$ .
- Calculate the magnetic susceptibility in zero magnetic field using the fluctuation-dissipation relation of the form

$$\frac{\chi(T)}{N} = \frac{1}{k_B T} \sum_{j=-N/2}^{N/2} \Gamma_{0j}, \quad (2)$$

in the thermodynamic limit,  $N \rightarrow \infty$ . For simplicity we assume  $N$  to be even. Note that  $\chi(T)$  is defined to be extensive, such that we obtain the intensive quantity by normalization with  $N$ , for details consider sections (3.4.5) and (3.4.6) of the lecture notes.

**Exercise 7.2 The Ideal Paramagnetic Gas and the Law of Mass Action**

The goal of this exercise is to understand the statistical mechanics of a mixture of ideal gases undergoing chemical reactions. The concepts will then be applied to an ideal gas of paramagnetic atoms which may combine to dimers whose magnetic moment vanishes.

Consider a gaseous mixture of  $r$  different substances  $A_1, \dots, A_r$  that undergo  $s$  chemical reactions,

$$\nu_1^\alpha A_1 + \dots + \nu_r^\alpha A_r = 0,$$

indexed by  $\alpha \in \{1, \dots, s\}$ , with  $\{\nu_i^\alpha\}$  the *stoichiometric coefficients* of the reaction  $\alpha$ .

*Example:* The reaction of water formation would be expressed through,  $A_1 = \text{H}_2$ ,  $A_2 = \text{O}_2$ , and  $A_3 = \text{H}_2\text{O}$ , with  $\nu_1 = 2$ ,  $\nu_2 = 1$ , and  $\nu_3 = -2$ .

Let  $N_i$  be the number of particles of the substance  $A_i$ . For a materially closed system, the set of possible variations in the number of particles is given by

$$dN_i = \sum_{\alpha=1}^s \nu_i^\alpha d\lambda^\alpha,$$

with independent variations  $d\lambda^1, \dots, d\lambda^s$  of the particle numbers according to the  $s$  reaction processes.

- a) Show that, assuming constant temperature and pressure, the condition for thermodynamic equilibrium constrains the chemical potentials  $\mu_i$  of the  $r$  species to obey

$$\sum_{i=1}^r \nu_i^\alpha \mu_i = 0,$$

for each  $\alpha = 1, \dots, s$  independently, and interpret this result.

- b) Let each substance  $A_i$  be an ideal gas composed of particles of mass  $m_i$  and with potential energy  $E_i$ . The Hamiltonian for the particles of type  $A_i$  then reads

$$\mathcal{H}_i = \sum_{j=1}^{N_i} \left( \frac{\vec{p}_j^2}{2m_i} + E_i \right).$$

Compute the grand canonical partition function  $\mathcal{Z}$  of the system and show that in equilibrium a *law of mass action*

$$\prod_{i=1}^r \langle N_i \rangle^{\nu_i^\alpha} = \prod_{i=1}^r \left[ V (2\pi m_i k_B T)^{3/2} e^{-\beta E_i} \right]^{\nu_i^\alpha} \equiv K^\alpha(T, V),$$

is obtained for each reaction  $\alpha$ .

This law states that in thermodynamic equilibrium, every chemical reaction is characterized by one value  $K^\alpha(T, V)$ , which depends only on the coefficients  $\nu_i^\alpha$ , the binding energies  $E_i$  and the particle masses  $m_i$ . In particular,  $K^\alpha(T, V)$  is independent of the proportion of species in the mixture.

*Remark:*  $E_i$  can be a binding energy of a molecule or a Zeeman energy for substances with atomic or molecular magnetic moments

- c) As a concrete example, consider an ideal paramagnetic gas in an external magnetic field  $H$ . Two particle species  $A_+$  ( $A_-$ ) of mass  $m$  have a magnetic moment  $|M|$  aligned parallel (antiparallel) to the field (with corresponding energy  $E = -MH$ ).

Two reactions are considered here. On the one hand,  $A_+$  and an  $A_-$  may combine to form a single molecule  $A_0$  whose magnetic moment vanishes. The energy released in this reaction is  $E_b$ . On the other hand we allow for a spin flip  $A_\sigma \mapsto A_{-\sigma}$ . Use the above results to compute average particle number of each species and the relative magnetization per particle

$$\sigma = M \frac{\langle N_+ \rangle - \langle N_- \rangle}{\langle N_+ + N_- + 2N_0 \rangle}.$$

Discuss the high and low temperature limits. Express the laws of mass action explicitly for these reactions.

**Office Hours:** Monday, November 7, 9-10 am (Roland Willa, HIT K 23.3)