

Exercise 12.1 Gaussian Fluctuations in the Ginzburg-Landau Model

Consider the Ginzburg-Landau model of the d -dimensional Ising model in presence of a magnetic field $H(\vec{r})$, introduced in chapter 5.3 of the lecture notes. Here, we only consider temperatures above the critical temperature T_c . In order to make the model exactly tractable, we assume that quartic fluctuations are negligible and ignore them. Therefore, the free energy functional for a given magnetization m and temperature T in d dimensions is given by

$$F(T, m, H) = \frac{1}{2} \int d^d r \left\{ A m(\vec{r})^2 - H(\vec{r}) m(\vec{r}) + \kappa [\vec{\nabla} m(\vec{r})]^2 \right\}, \quad (1)$$

where $A = a\tau$, with $\tau = (T - T_c)/T_c$. For the calculations we assume our system to be a cube of side length L with periodic boundary conditions on m .

- a) Use the Fourier transform of the magnetization field,

$$m(\vec{r}) = \frac{1}{\sqrt{L^d}} \sum_{\vec{q}} m_{\vec{q}} e^{i\vec{q}\cdot\vec{r}}, \quad (2)$$

and compute the energy functional $F(T, m)$ in the transformed coordinates $\{m_{\vec{q}}\}$. Which values of \vec{q} are allowed in the sum and which values of \vec{q} are independent? Note that $m(\vec{r})$ is real and interpret its implication on the $m_{\vec{q}}$.

- b) Compute the canonical partition function,

$$Z(T) = \int \mathcal{D}m e^{-F(T, m)/k_B T}, \quad (3)$$

and the free energy $F(T) = -k_B T \log Z(T)$ by using Gaussian integration. Argue that the finite number of degrees of freedom (finite lattice spacing) of our Ising model introduces a momentum cutoff Λ , which is crucial to regulate the otherwise ill defined integrals (cf. Debye wave vector for phonons).

Hint: Rewrite the functional measure $\mathcal{D}m$ according to

$$\mathcal{D}m = \prod_{\vec{q}} dm_{\vec{q}} dm_{-\vec{q}}. \quad (4)$$

Why do we use $dm_{\vec{q}} dm_{-\vec{q}}$?

- c) Compute the internal energy and the specific heat c_V in the thermodynamic limit $L \rightarrow \infty$ for vanishing external field ($H(\vec{r}) \equiv 0$). Study its behavior near the critical temperature where $\tau = 0$. Compare the critical exponent of c_V with the mean field result of section 5.3.2 for different dimensions d .
- d) Derive an expression for the magnetic susceptibility, defined as the negative second derivative of the free energy with respect to the external field H in the limit of vanishing field, i.e.

$$\chi(T) = - \left. \frac{\partial^2 F(T)}{\partial H^2} \right|_{H=0}. \quad (5)$$

What is the critical exponent of χ ? Compare the result with the mean field result of section 5.2.5 of the lecture notes, Eq. (5.58).