

# Solutions 1 - reminder of Newtonian mechanics

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## 1. Geostationary space station

a) The observer must be on the equator of the earth. The orbit of the space station is a large circle in the equatorial plane with center at the center of the earth.

b) Since the centripetal force is equal to the gravitational force we have

$$\frac{mv^2}{R} = \frac{GM_{\oplus}m}{R^2}, \quad (1)$$

where  $M_{\oplus}$  is the mass of the earth,  $v$  is the circular velocity and  $R$  the orbital radius of the space station. With

$$v = \frac{s}{t} = \frac{2\pi R}{T} \quad (2)$$

we find the following distance  $L$  between the observer and the space station:

$$L = R - R_{\oplus} = \left( \frac{GMT^2}{4\pi^2} \right)^{1/3} - R_{\oplus}. \quad (3)$$

with  $T = 24h$ ,  $M_{\oplus} = 6.0 \cdot 10^{25}kg$  and  $R_{\oplus} = 6.4 \cdot 10^6m$  we finally obtain

$$L = 3.6 \cdot 10^4 km. \quad (4)$$

## 2. Circling particle

a) The tension in the string provides the centripetal force needed for the circular motion, hence  $F = mv_0^2/R_0$ .

b) The angular momentum of the mass  $m$  is  $J = mv_0R_0$ , the kinetic energy is  $T = mv_0^2/2$ .

c) The radius of the circular motion of the mass  $m$  decreases when the tension in the string is increased gradually. The angular momentum is conserved, thus

$$mv_0R_0 = mv_1\left(\frac{R_0}{2}\right) \quad \implies \quad v_1 = 2v_0. \quad (5)$$

The final kinetic energy is then

$$T_1 = \frac{1}{2}mv_1^2 = 2mv_0^2. \quad (6)$$

The reason why the pulling of the string should be gradual is that the vectors  $\mathbf{r}$  and  $\mathbf{v}$  must stay approximately perpendicular.

### 3. Fast rotating planet

The conservation of energy can be written as the statement

$$E = T(R) + U(R) = T(\infty) + U(\infty) = \text{const.} \quad (7)$$

The escape velocity  $v_e$  can now be determined by setting  $T(\infty)$  to zero. This leads to

$$\begin{aligned} \frac{1}{2}mv_e^2(R) &= \int_R^\infty F(r)dr = GMm \int_R^\infty \frac{1}{r^2}dr = \frac{GMm}{R}, \\ v_e(R) &= \sqrt{\frac{2GM}{R}}. \end{aligned} \quad (8)$$

On our fast rotation planet we have  $g_{pol} = 2g_{eq}$ , as well as

$$g_{pol} = \frac{GM}{R^2} \quad g_{eq} = \frac{GM}{R^2} - \frac{v^2}{R}, \quad (9)$$

what leads to the equation

$$\frac{GM}{R} = 2v^2. \quad (10)$$

Substituting eq. (10) into eq. (8) gives the final result

$$v_e = 2v. \quad (11)$$

### 3. Sphere

a) Conservation of energy gives (see Fig. 1)

$$\begin{aligned} E &= mg(2R) = \frac{1}{2}mv^2 + mgR(1 + \cos\theta), \\ \frac{1}{2}mv^2 &= mgR(1 - \cos\theta). \end{aligned} \quad (12)$$

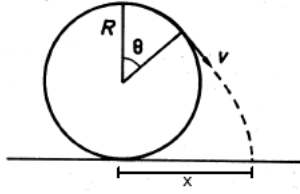


Figure 1: Particle falling from a sphere.

The radial force the sphere exerts on the particle is

$$F = mg \cos \theta - \frac{mv^2}{R}. \quad (13)$$

When  $F = 0$ , the constraint vanishes and the particle leaves the sphere. At this instance we have

$$v^2 = gR \cos \theta \quad v^2 = 2gR(1 - \cos \theta) \quad (14)$$

giving

$$\cos \theta = 2/3, \quad v = \sqrt{\frac{2gR}{3}}. \quad (15)$$

b) After leaving the sphere, the particle follows a parabolic trajectory until it hits the plane. The trajectory is described by

$$\begin{aligned} x(t) &= vt \cos \theta + x_0 \\ y(t) &= -\frac{1}{2}gt^2 - vt \sin \theta + y_0 \end{aligned} \quad (16)$$

with the initial conditions  $x_0 = R \sin \theta$  and  $y_0 = R(1 + \cos \theta)$ . This can also be written as

$$y(x) = y_0 - \tan \theta (x - x_0) - \frac{g}{2v^2 \cos^2 \theta} (x - x_0)^2 \quad (17)$$

Setting  $y(x) = 0$  and solving for  $x$  leads to

$$x - x_0 = \frac{v}{g} \cos \theta \left( -v \sin \theta \pm \sqrt{v^2 \sin^2 \theta + \frac{1}{2}gy_0} \right). \quad (18)$$

Substituting  $\theta$  and  $v$  and only considering the positive solution leads to

$$x = \frac{\sqrt{5}}{27}(5 + 2\sqrt{13})R \simeq 1.01R. \quad (19)$$

The result is independent of  $g$ !

## 5. Enjoying wine in a train

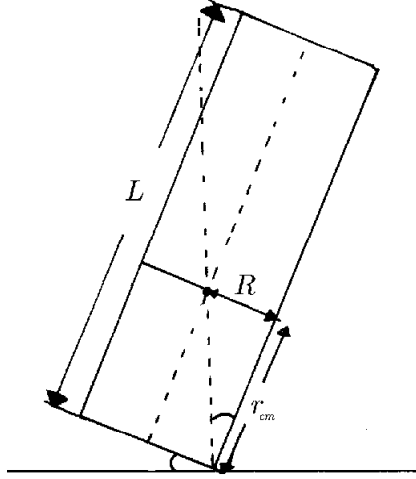


Figure 2: Illustration of the the bottle.

Let us study a bottle that has been tilted as in Fig. 2. To simplify calculations, we ignore the fact that the surface of the wine is always horizontal. When the center of mass  $r_{\text{cm}}$  lies on the  $y$ -axis, there is no net torque. Let us call the angle that gives us this configuration  $\theta_{\text{eq}}$ . For  $\theta < \theta_{\text{eq}}$ , torque will force the bottle back into a standing position and for  $\theta > \theta_{\text{eq}}$  the bottle will fall. In essence, the larger  $\theta_{\text{eq}}$  is, the more stable is the bottle!

The setup is shown in Fig. 2. The equilibrium angle is calculated through

$$R/r_{\text{cm}} = \tan(\theta_{\text{eq}}), \quad (20)$$

where  $R$  is the radius of the bottle. Maximizing  $\theta_{\text{eq}}$  is equivalent to minimizing  $r_{\text{cm}}$  which is given by

$$r_{\text{cm}} = \frac{m_b \frac{L}{2} + m_w x \frac{L}{2}}{m_b + m_w}, \quad (21)$$

where  $m_b$  is the mass of an empty wine bottle,  $m_w$  is mass of the wine and  $x$  is the fraction of wine in the bottle. By using the fact that the wine mass is  $m_w = x m_{w,f}$ , where  $m_{w,f}$  is the mass of wine in a full bottle we can expand into

$$r_{\text{cm}} = \frac{m_b \frac{L}{2} + m_{w,f} \frac{L}{2} x^2}{m_b + m_{w,f} x}. \quad (22)$$

The sought solution is when  $\frac{dr_{\text{cm}}}{dx} = 0$  and  $\frac{d^2 r_{\text{cm}}}{dx^2} > 0$  (minimum of function).

The derivative is

$$\frac{dr_{\text{cm}}}{dx} = \frac{(m_b + m_{w,f}x)m_{w,f}Lx - m_{w,f}(m_b \frac{L}{2} + m_{w,f} \frac{L}{2}x^2)}{(m_b + m_{w,f}x)^2} \quad (23)$$

Equating this to zero and simplifying renders the quadratic equation

$$x^2 + 2 \frac{m_b}{m_{w,f}}x - \frac{m_b}{m_{w,f}} = 0 \quad (24)$$

which has the solution

$$x = -\frac{m_b}{m_{w,f}} \pm \sqrt{\left(\frac{m_b}{m_{w,f}}\right)^2 + \frac{m_b}{m_{w,f}}}. \quad (25)$$

We see that the optimum amount of wine in the bottle only depends on the mass of the glass bottle itself and the mass of the wine itself *in a full bottle*. By using  $m_b = 450g$  and  $m_{w,f} = 750g$  we get  $x = 0.3798 \approx 3/8$ . For a standard bottle of wine,  $R \approx L/6$ . Using this, and the solution for  $x$ , in Eq. 20 we see that the maximum angle of stability is  $\theta_{\text{eq}} \approx 23.7$ .