

## Exercises for "Phenomenology of Particle Physics I"

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### Exercise 18

Show the following distribution identity:

$$\int d^3x f_p(x)^* i \overleftrightarrow{\partial}_0 f_q(x) \Big|_{t=0, p^0=E_{\vec{p}}, q^0=E_{\vec{q}}} = \delta^3(\vec{p} - \vec{q})$$

where we define  $f(x) \overleftrightarrow{\partial}_x g(x) := -g(x) \partial_x f(x) + f(x) \partial_x g(x)$  and we have  $f_p(x) = \frac{1}{\sqrt{(2\pi)^3}} \frac{1}{\sqrt{2E_{\vec{p}}}} e^{-ipx}$ .

### Exercise 19

Starting from the fourier transform of the real Klein-Gordon field and the canonical momentum density conjugate to it (in the Heisenberg picture)

$$\phi(x) = \int d^3p \frac{1}{\sqrt{(2\pi)^3}} \frac{1}{\sqrt{2E_{\vec{p}}}} \left( f_p(x) a(\vec{p}) + f_p(x)^* a(\vec{p})^\dagger \right) \Big|_{p^0=E_{\vec{p}}}$$

$$\Pi(x) = \partial_0 \phi(x) = \int d^3p \frac{1}{\sqrt{(2\pi)^3}} \frac{1}{\sqrt{2E_{\vec{p}}}} (iE_{\vec{p}}) \left( -f_p(x) a(\vec{p}) + f_p(x)^* a(\vec{p})^\dagger \right) \Big|_{p^0=E_{\vec{p}}}$$

with  $E_{\vec{p}} = (|\vec{p}|^2 + m^2)^{1/2}$ , derive the expressions for  $a(\vec{p})$  and  $a(\vec{p})^\dagger$  by inversion. The result is

$$a(\vec{p}) = \sqrt{(2\pi)^3} \sqrt{2E_{\vec{p}}} \int d^3x f_p(x)^* i \overleftrightarrow{\partial}_0 \phi(x)$$

$$a(\vec{p})^\dagger = \sqrt{(2\pi)^3} \sqrt{2E_{\vec{p}}} \int d^3x \phi(x) i \overleftrightarrow{\partial}_0 f_p(x).$$

### Exercise 20

Show that if we postulate the commutation relations

$$\begin{aligned} [a(\vec{p}), a(\vec{q})^\dagger] &= \delta^3(\vec{p} - \vec{q})(2\pi)^3 2E_{\vec{p}} \\ [a(\vec{p}), a(\vec{q})] &= [a(\vec{p})^\dagger, a(\vec{q})^\dagger] = 0 \end{aligned}$$

for  $a$  and  $a^\dagger$  we arrive at the following commutation relations for the field and the canonical momentum density conjugate to it

$$\begin{aligned} [\phi(\vec{x}, t), \Pi(\vec{x}', t)] &= i\delta^3(\vec{x} - \vec{x}') \\ [\phi(\vec{x}, t), \phi(\vec{x}', t)] &= [\Pi(\vec{x}, t), \Pi(\vec{x}', t)] = 0. \end{aligned}$$

### Exercise 21 (corrected)

We define

$$\Delta^\pm(x) = - \int_{C^\pm} \frac{dp^0}{2\pi} \int \frac{d^3p}{(2\pi)^3} \frac{e^{-ipx}}{p^2 - m^2}$$

where  $C^+$  and  $C^-$  are contours in the complex  $p^0$ -plane,  $C^+$  goes around  $p^0 = E_{\vec{p}}$  once in counterclockwise direction,  $C^-$  goes around  $p^0 = -E_{\vec{p}}$  once in counterclockwise direction.

- Use the residue theorem and the formula

$$\int \frac{d^3p}{2E_{\vec{p}}} = \int d^4p \delta(p^2 - m^2) \Theta(p^0)$$

to show

$$\Delta^\pm(x) = \mp i \int \frac{d^4p}{(2\pi)^3} \delta(p^2 - m^2) e^{\mp ipx} \Theta(p^0).$$

- Show that  $[\phi(x), \phi(y)] = i\Delta(x - y) = i(\Delta^+(x - y) + \Delta^-(x - y))$  vanishes for spacelike  $((x - y)^2 < 0)$  separation  $x - y$ .