

Exercise 1. Renormalization in QCD at one loop: The Vertices

In this series we will renormalise the vertices of QCD at one-loop. Here you find all the results of the renormalisation in the \overline{MS} . Try to do some of these integrals by your own (you don't have to compute all of them, but try to become confident with these kind of calculations).

We define the renormalised quantities as

$$\begin{aligned} A_{\mu,0}^a &= \sqrt{Z_A} A_{\mu,R}^a & g_0 &= Z_g g_R \\ \eta_{a,0} &= \sqrt{\tilde{Z}_3} \eta_{a,R} & m_0 &= Z_m m_R \\ \psi_0 &= \sqrt{Z_2} \psi_R & \xi_0 &= Z_A \xi_R \end{aligned}$$

(a) Ghost-gluon vertex

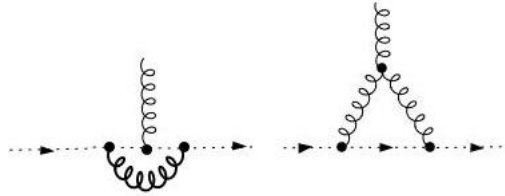


Figure 1: One loop ghost-gluon vertexes

$$I = -i g p_\mu (-i f_{abc}) \left(\tilde{Z}_1 + \frac{\alpha_s}{4\pi} N_C (4\pi)^\epsilon \Gamma(\epsilon) + finite \right) \tag{1}$$

$$\begin{aligned} \tilde{Z}_1 &= Z_g \tilde{Z}_3 Z_A^{1/2} \\ &= 1 - \frac{\alpha_s}{4\pi} N_C (4\pi)^\epsilon \Gamma(\epsilon) \end{aligned} \tag{2}$$

(b) Quark-gluon vertex

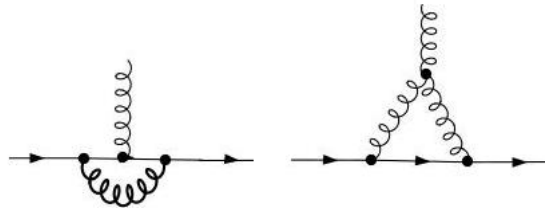


Figure 2: One loop quark-gluon vertexes

$$I = -i g \gamma_\mu t_{ij}^a \left(Z_{1F} + \frac{\alpha_s}{4\pi} (N_C + C_F) (4\pi)^\epsilon \Gamma(\epsilon) + finite \right) \tag{3}$$

$$\begin{aligned} Z_{1F} &= Z_g Z_2 Z_A^{1/2} \\ &= 1 - \frac{\alpha_s}{4\pi} (N_C + C_F) (4\pi)^\epsilon \Gamma(\epsilon) \end{aligned} \tag{4}$$

(c) 3 gluon-gluon vertex

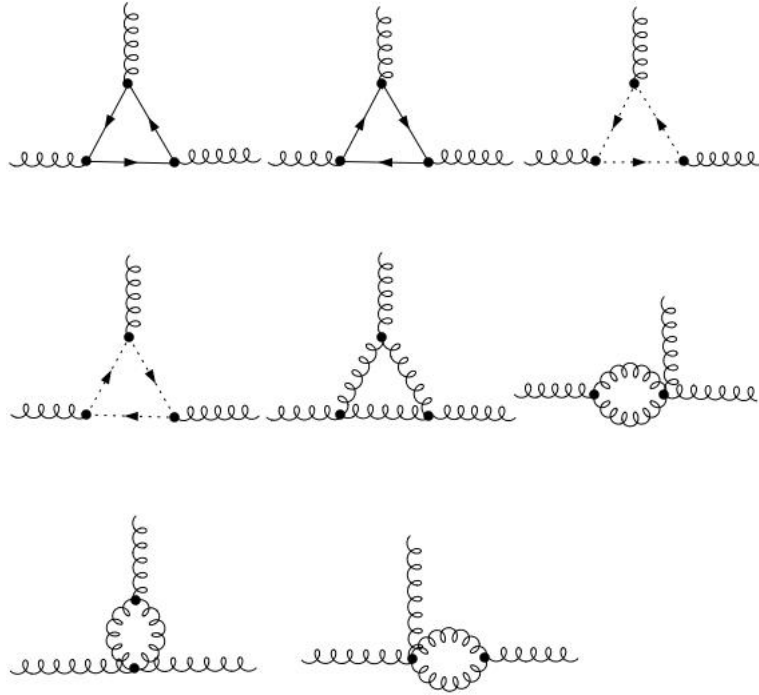


Figure 3: One loop 3 gluon-gluon vertexes

$$\begin{aligned}
 I &= -i g f^{abc} V_{\mu\nu\sigma}(p_1, p_2, p_3) \left(Z_1 + \frac{\alpha_s}{4\pi} \left(N_C \left(-\frac{2}{3} \right) + \frac{4}{3} \frac{1}{2} n_f \right) (4\pi)^\varepsilon \Gamma(\varepsilon) + finite \right) \\
 Z_1 &= Z_g Z_A^{3/2} \\
 &= 1 - \frac{\alpha_s}{4\pi} \left(N_C \left(-\frac{2}{3} \right) + \frac{4}{3} \frac{1}{2} n_f \right) (4\pi)^\varepsilon \Gamma(\varepsilon)
 \end{aligned} \tag{6}$$

(d) 4 gluon-gluon vertex

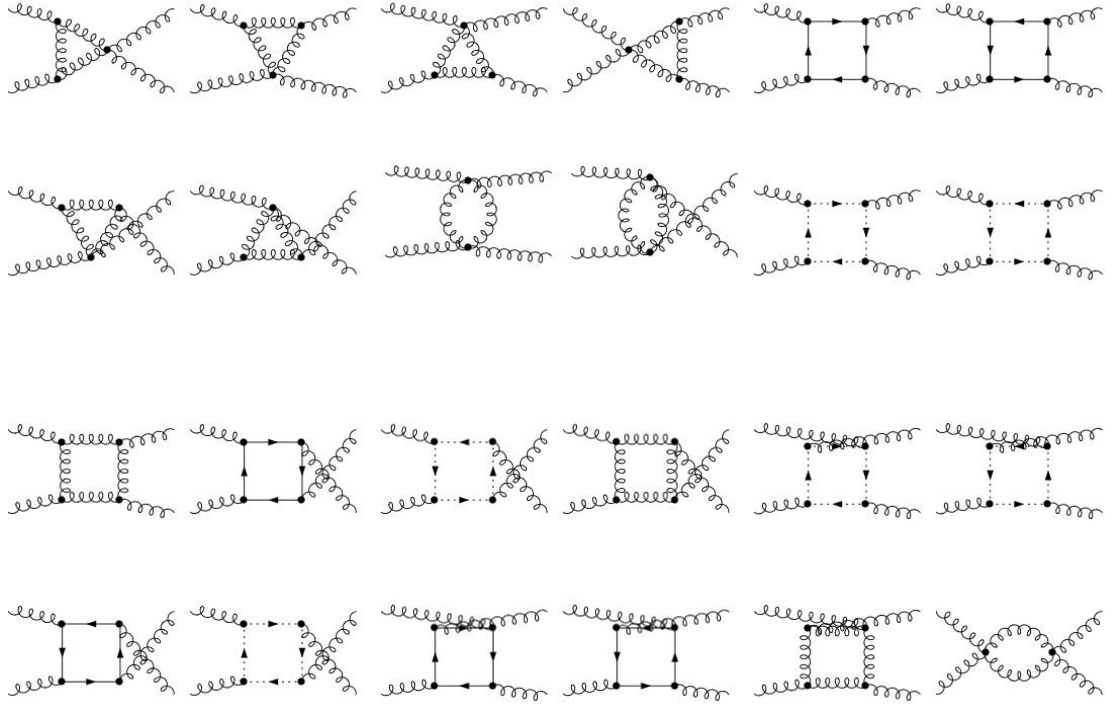


Figure 4: One loop 4 gluon-gluon vertexes

$$I = -i g^2 V_{\mu\nu\rho\sigma}^{abcd} \left(Z_4 + \frac{\alpha_s}{4\pi} \left(\frac{N_C}{3} + \frac{4}{3} \frac{1}{2} n_f \right) (4\pi)^\varepsilon \Gamma(\varepsilon) + \text{finite} \right) \quad (7)$$

$$\begin{aligned} Z_4 &= Z_g^2 Z_A^2 \\ &= 1 - \frac{\alpha_s}{4\pi} \left(\frac{N_C}{3} + \frac{4}{3} \frac{1}{2} n_f \right) (4\pi)^\varepsilon \Gamma(\varepsilon) \end{aligned} \quad (8)$$

Now define $\Delta_\varepsilon \equiv \frac{\alpha_s}{4\pi} (4\pi)^\varepsilon \Gamma(\varepsilon)$. We have

$$Z_A = 1 - \left[\frac{2}{3} n_f - \frac{5}{3} N_C \right] \Delta_\varepsilon \quad (9)$$

$$\tilde{Z}_1 + \frac{N_C}{2} \Delta_\varepsilon \quad (10)$$

$$Z_2 = 1 - C_F \Delta_\varepsilon \quad (11)$$

$$Z_m = 1 - 3C_F \Delta_\varepsilon \quad (12)$$

$$Z_1 = 1 - \frac{\alpha_s}{4\pi} \left(N_C \left(-\frac{2}{3} \right) + \frac{4}{3} \frac{1}{2} n_f \right) \Delta_\varepsilon \quad (13)$$

$$\tilde{Z}_1 = 1 - \frac{N_C}{2} \Delta_\varepsilon \quad (14)$$

$$Z_{1F} = 1 - \frac{\alpha_s}{4\pi} (N_C + C_F) \Delta_\varepsilon \quad (15)$$

$$Z_4 = 1 - \frac{\alpha_s}{4\pi} \left(\frac{N_C}{3} + \frac{4}{3} \frac{1}{2} n_f \right) \Delta_\varepsilon \quad (16)$$

$$\frac{Z_1}{Z_A} = \frac{\tilde{Z}_1}{\tilde{Z}_3} = \frac{Z_{1F}}{Z_2} = \frac{Z_4}{Z_1} = Z_g Z_A^{1/2} = 1 - N_C \Delta_\varepsilon \quad (17)$$

Exercise 2. Colour Factors

It is common to adopt a normalisation convention based on the trace of products for the matrices t^a belonging to an irreducible representation. In particular, for the $SU(N_C)$ Lie Algebra, one usually uses

$$Tr(t^a t^b) = T_F \delta^{ab}, \quad T_F = \frac{1}{2}$$

The matrices t^a are hermitian, therefore there are $N_C^2 - 1$ generators which form a basis for this space of hermitian matrices. We then have that their commutators are i times hermitian traceless matrices, i.e.

$$[t^a, t^b] = i f^{abc} t^c, \quad i f^{abc} = \frac{1}{T_F} Tr([t^a, t^b] t^c)$$

where f^{abc} are real constants. We also know that actually they are the generators in the adjoint representation:

$$(t^c)^{ab} = i f^{acb}$$

In order to calculate the colour factors, we know that

$$\begin{aligned} Tr(\text{fermion loop}) &= N_C \\ Tr(t^a) &= 0 \\ Tr(t^a t^b) &= T_F \delta^{ab} \\ t^a t^a &= C_F, \quad C_F = \frac{1}{2} \frac{N_C^2 - 1}{N_C} \end{aligned} \quad (18)$$

Show that this other useful identity holds

$$(t^a)_j^i (t^a)_l^k = T_F \left[\delta_l^i \delta_j^k - \frac{1}{N_C} \delta_j^i \delta_l^k \right] \quad (19)$$

Hint. First notice that the general form of this expression is

$$(t^a)_j^i (t^a)_l^k = a [\delta_l^i \delta_j^k - b \delta_j^i \delta_l^k]$$

Then, multiply it once by δ_i^j and the other time by $(t^b)_i^j$ in order to find who b and a are.