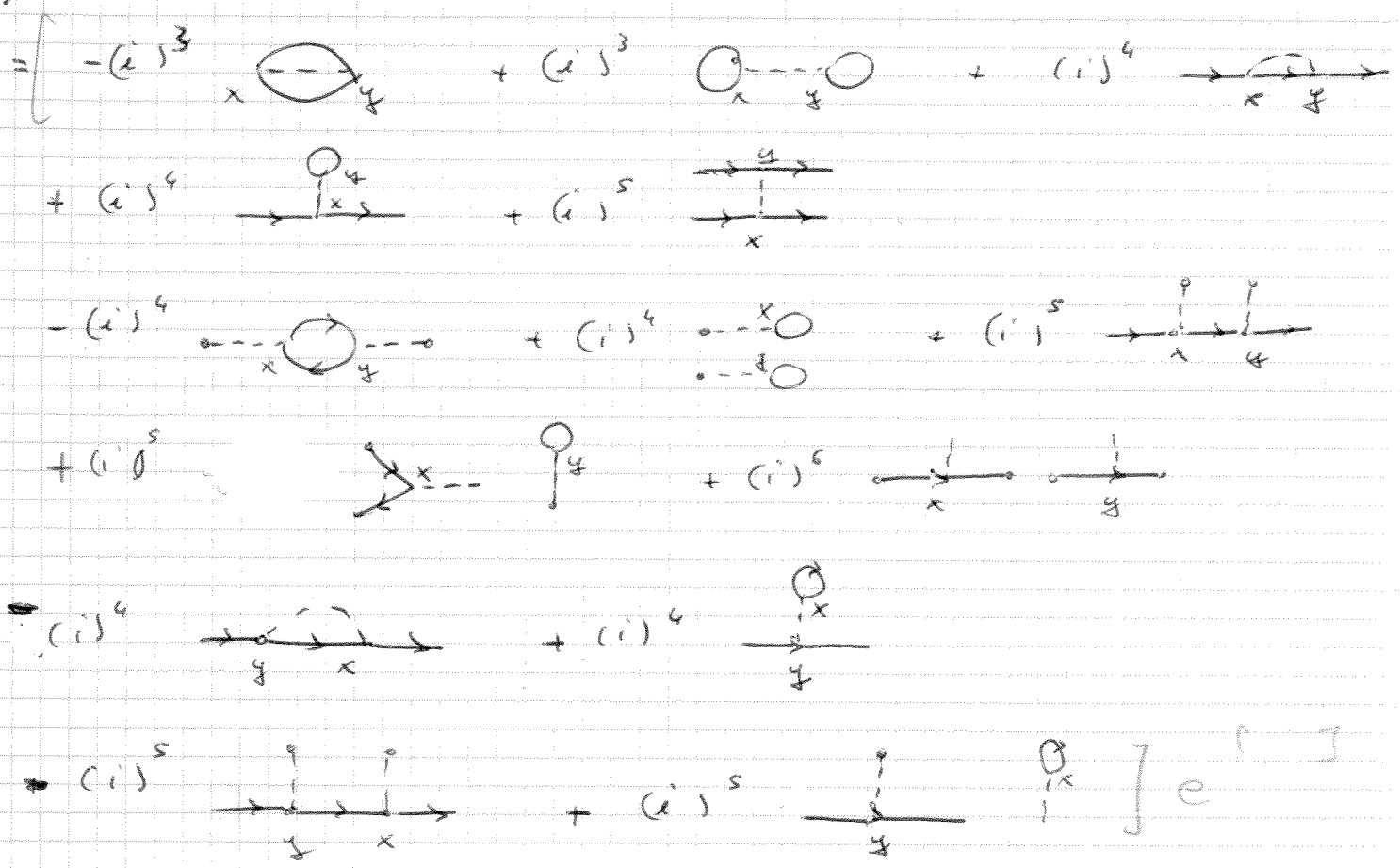


$$\begin{aligned} \overline{Z}[\gamma, \bar{\gamma}, J] &= N \int \Delta\bar{\psi} \Delta\psi \Delta\phi e^{iS_0 + iS_{se} + i\int d^4x (\lambda\bar{\psi}\psi\phi + J\psi + \bar{\psi}J\psi)} \\ &= N e^{iS_{int}} Z_0[\gamma, \bar{\gamma}, J] \end{aligned} \quad (1)$$

$$\begin{aligned} Z[\gamma, \bar{\gamma}, J] &= N \left[1 + i\lambda \int d^4x \frac{\delta}{\delta\psi(x)} \frac{\delta}{\delta\bar{\psi}(x)} \frac{\delta}{\delta J(x)} + \right. \\ &\quad \left. \frac{(i\lambda)^2}{2} \int d^4x \int d^4y \frac{\delta}{\delta\psi(x)} \frac{\delta}{\delta\bar{\psi}(y)} \frac{\delta}{\delta J(y)} \frac{\delta}{\delta\psi(y)} \frac{\delta}{\delta\bar{\psi}(x)} \frac{\delta}{\delta J(x)} \right] e^{\dots} \end{aligned}$$



on

Normalization : vacuum diag.

$$Z[0,0,0] = 1 = N \int \Delta\bar{\psi} \Delta\psi \Delta\phi e^{iS_0 + iS_{sc} + i\lambda \int d^4x \psi^4}$$

$$\Rightarrow \frac{1}{N} = e^{i\lambda \int d^4x \frac{iS}{S\psi(x)} \frac{S}{iS\bar{\psi}(x)} \frac{S}{iS\phi(x)} Z_0[\bar{\psi}, \psi, \phi] \Big|_{S=0}} = 1 + \frac{(i\lambda)^2}{2!} [-i^3 \text{loop} + i^3 \text{tadpole}] + o(\lambda^3)$$

Scalar prop.

$$\frac{S^2}{iS(z_1) iS(z_2)} \quad | \quad iW[z] \quad | \quad \bar{\psi} = \psi = 0$$

$$= \frac{S}{iS(z_1)} \frac{1}{z[z]} \frac{S}{iS(z_2)} \quad | \quad z[z] \quad | \quad \psi = 0$$

$$= \frac{1}{(i)^2} \left[- \frac{1}{z^2[z]} \frac{S}{iS(z_1)} \frac{S}{iS(z_2)} \quad | \quad z[z] \quad | \quad \psi = 0 \right. \\ \left. + \frac{1}{z[z]} \frac{S^2}{iS(z_1) iS(z_2)} \quad | \quad z[z] \quad | \quad \psi = 0 \right]$$

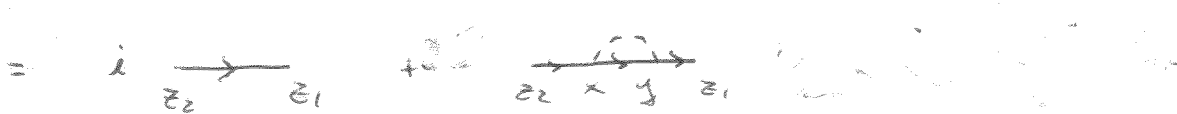


Fermion prop.

$$\frac{S}{S\psi(z_1)} \frac{S}{iS\bar{\psi}(z_2)} \quad | \quad iW[z] \quad | \quad \bar{\psi} = \psi = 0$$

$$= \frac{S}{S\psi(z_1)} \frac{1}{z[z]} \frac{S}{iS\bar{\psi}(z_2)} \quad | \quad z[z] \quad | \quad = 0$$

$$= - \frac{1}{z^2[z]} \frac{S}{S\psi(z_1)} \frac{S}{iS\bar{\psi}(z_2)} \quad | \quad z[z] \quad | \quad + \frac{1}{z[z]} \frac{S^2}{S\psi(z_1) iS\bar{\psi}(z_2)} \quad | \quad z[z] \quad | \quad = 0$$



Scattering 4 fermions:

(3)

$$\frac{S}{S_4(z_4)} \frac{S}{S_4(z_3)} \frac{S}{S_4(z_2)} \frac{S}{S_4(z_1)} \quad iW \quad \Big|_{S=0}$$

$$= i \left(\text{diagram 1} - \text{diagram 2} \right)$$

Scattering $2f \rightarrow 2sc$

$$\frac{1}{(i)^2} \frac{S}{S_4} \frac{S}{S_4} \frac{S}{S_4} \frac{S}{S_4} \quad iW \quad \Big|_{S=0}$$

$$= -i \left(\text{diagram 1} + \text{diagram 2} \right)$$

ex. 2

(4)

Under a global transformation, we have

$$\psi' = e^{i\alpha} \psi$$

$$j^\mu = \frac{\delta \mathcal{L}}{\delta(\partial_\mu \psi)} \delta \psi = \bar{\psi} \gamma^\mu \psi$$

So that the conserved charge is

$$Q = \int d^3x j^0(x) = \int d^3x \bar{\psi} \gamma^0 \psi = \int d^3x \psi^\dagger \psi$$

ex. 3

a) Infinitesimally we have

$$U(\theta) = 1 + ig \theta_a T^a$$

$$U(\theta)^{-1} = 1 - ig \theta_a (T^a)^\dagger$$

From Unitarity we get

$$U^\dagger U = 1 \approx (1 - ig (T^a)^\dagger \theta_a) (1 + ig T^b \theta_b) \approx 1 + \mathcal{O}(\theta^2)$$

so that $(T^a)^\dagger = T^a$

Computing $(U^b)^{-1} U^{-1} U^b U =$

$$= (1 - ig T^a \theta_a - \frac{1}{2} g^2 (T^a)^2 (\theta^a)^2) (1 - ig T^b \theta_b - \frac{1}{2} g^2 (T^b)^2 \theta_b^2) \\ \cdot (1 + ig T^a \theta_a - \frac{1}{2} g^2 (T^a)^2 (\theta^a)^2) (1 + ig T^b \theta_b - \frac{1}{2} g^2 (T^b)^2 \theta_b^2)$$

$$= 1 - g^2 [T^a, T^b] \theta^a \theta^b + \mathcal{O}(\theta^3)$$

Since $U, U' \in SU(N)$ are chosen to be independent elements, this result must be a non-trivial element of $SU(N)$ for $a \neq b$. Therefore, we can write

$$= 1 + i g^2 f^{abc} \theta^a \theta^b T^c + O(\theta^3)$$

b) To show that f^{abc} are real, consider

$$([T^a, T^b])^\dagger = -i (f^{abc})^\dagger (T^c)^\dagger$$

$$- [T^a, T^b] = -i (f^{abc})^\dagger T^c$$

$$[T^a, T^b] = i (f^{abc})^\dagger T^c$$

\Downarrow

$$(f^{abc})^\dagger = f^{abc}$$