

Exercise 1. *Electron in a time-dependent magnetic field*

Imagine an electron at rest at the origin, in the presence of a magnetic field whose magnitude is constant, but whose direction sweeps out a cone, of opening angle α , at a constant angular velocity ω :

$$\vec{B}(t) = B_0 \left[\sin \alpha \cos \omega t \hat{i} + \sin \alpha \sin \omega t \hat{j} + \cos \alpha \hat{k} \right]. \quad (1)$$

The Hamiltonian is

$$H(t) = \frac{e}{m} \vec{B} \cdot \vec{S} = \frac{\hbar \omega_1}{2} \begin{pmatrix} \cos \alpha & e^{-i\omega t} \sin \alpha \\ e^{i\omega t} \sin \alpha & -\cos \alpha \end{pmatrix}, \quad (2)$$

where $\omega_1 = \frac{eB_0}{m}$.

- (a) Compute the normalized eigenspinors of $H(t)$:

$$H(t)\chi_{\pm}(t) = E_{\pm}\chi_{\pm},$$

and show that

$$E_{\pm} = \pm \frac{\hbar \omega_1}{2}.$$

- (b) Supposing that the electron starts out with spin up along $\vec{B}(0)$:

$$\chi(0) = \begin{pmatrix} \cos(\alpha/2) \\ \sin(\alpha/2) \end{pmatrix},$$

verify that the exact solution of the time-dependent Schroedinger equation can be written as:

$$\begin{aligned} \chi(t) = & \left[\cos\left(\frac{\lambda t}{2}\right) - i \frac{\omega_1 - \omega \cos \alpha}{\lambda} \sin\left(\frac{\lambda t}{2}\right) \right] e^{-i\omega t/2} \chi_+(t) \\ & + i \left[\frac{\omega}{\lambda} \sin \alpha \sin\left(\frac{\lambda t}{2}\right) \right] e^{+i\omega t/2} \chi_-(t). \end{aligned} \quad (3)$$

- (c) Compute the probability of a transition to spin down exactly.

Consider then the two regimes:

- $\omega \ll \omega_1$ (the so-called *adiabatic regime*)
- $\omega \gg \omega_1$

and compare the two behaviors. Is the adiabatic theorem satisfied?

- (d) If $\omega_1 \ll \omega$, $H(t)$ can be considered as a time-dependent perturbation. Assuming that the perturbation is switched on at $t > 0$, compute the first order perturbative correction to the eigenstate

$$\chi(0) = \begin{pmatrix} \cos(\alpha/2) \\ \sin(\alpha/2) \end{pmatrix}$$

in time-dependent perturbation theory and compare it to the expansion of the exact solution (3).

- (e) Consider now again the adiabatic regime $\omega \ll \omega_1$. Starting from the exact solution (3), show that in this regime it can be written as:

$$\chi(t) = e^{i(\epsilon(t) + \gamma(t))} \chi_+(t) + \mathcal{O}(\omega/\omega_1)$$

where $\epsilon(t)$ is the so-called dynamic phase, and $\gamma(t)$ is the geometric phase.

- (f) The geometric phase can also be computed directly from $\chi_+(t)$, using the adiabatic approximation. Compute the geometric phase and compare the result with what you obtained in point (e).