

Exercise 1. Free Electron Gas.

The Hamiltonian of a gas of N free electrons is written in the second quantization formalism as

$$H = \sum_{\mathbf{k},s} \xi_{\mathbf{k}} c_{\mathbf{k},s}^{\dagger} c_{\mathbf{k},s} , \quad (1)$$

where $c_{\mathbf{k},s}$ (resp. $c_{\mathbf{k},s}^{\dagger}$) is the annihilation (resp. creation) operator of the electron mode \mathbf{k}, s of energy $\xi_{\mathbf{k}} = \epsilon_{\mathbf{k}} - \mu$. (Here $\epsilon_{\mathbf{k}} = \hbar^2 k^2 / (2m)$ and μ is the chemical potential, $\mu = E_F$ at $T = 0$.) The index s distinguishes the two spin components.

Let's look at excitations that are holes under the Fermi level and electrons above the Fermi level. We would like to rewrite the Hamiltonian in a form which involves explicitly only these excitations. We define the creation and annihilation operators of an excitation $\alpha_{\mathbf{k},s}^{\dagger}, \alpha_{\mathbf{k},s}$ by

$$\alpha_{\mathbf{k},\uparrow} = \begin{cases} c_{\mathbf{k},\uparrow} & \text{for } k > k_F \\ c_{-\mathbf{k},\downarrow}^{\dagger} & \text{for } k < k_F \end{cases} ; \quad \alpha_{\mathbf{k},\downarrow} = \begin{cases} c_{\mathbf{k},\downarrow} & \text{for } k > k_F \\ c_{-\mathbf{k},\uparrow}^{\dagger} & \text{for } k < k_F \end{cases} . \quad (2)$$

- (a) Show that the α, α^{\dagger} 's obey fermionic commutation relations.
- (b) Argue that eq. (2) is a unitary transformation of the creation and annihilation operators. Such a transformation is also called a *Bogoliubov transformation*. What happens if you act with the annihilators $c_{\mathbf{k},s}$ and $\alpha_{\mathbf{k},s}$ on the ground state of the gas?
- (c) Rewrite the Hamiltonian (1) in the form

$$H = \sum_{\mathbf{k}} |\xi_{\mathbf{k}}| \left(\alpha_{\mathbf{k}\uparrow}^{\dagger} \alpha_{\mathbf{k}\uparrow} + \alpha_{\mathbf{k}\downarrow}^{\dagger} \alpha_{\mathbf{k}\downarrow} \right) + E_G ; \quad E_G = 2 \sum_{k < k_F} \xi_{\mathbf{k}} . \quad (3)$$

Exercise 2. Correlation Functions in a Fermi Sea.

Consider a gas of N identical fermions with spin $1/2$. The fermions are free and non-interacting. The ground state is then given by

$$|\Phi_0\rangle = \prod_{|\mathbf{k}| \leq k_F, s} a_{\mathbf{k}s}^{\dagger} |0\rangle . \quad (4)$$

One defines the *one-particle correlation function* $G_s(\mathbf{x} - \mathbf{y})$ as

$$G_s(\mathbf{x} - \mathbf{y}) = \frac{n}{2} g_s(\mathbf{x} - \mathbf{y}) = \langle \Phi_0 | \Psi_s^{\dagger}(\mathbf{x}) \Psi_s(\mathbf{y}) | \Phi_0 \rangle . \quad (5)$$

This is the amplitude of recreating a fermion of spin s at position \mathbf{x} when one was annihilated at position \mathbf{y} with same spin.

- (a) Using explicit expressions for the field operators $\Psi_s(\mathbf{x})$, calculate $G_s(\mathbf{x} - \mathbf{y})$ and sketch its graph as function of $|\mathbf{x} - \mathbf{y}|$. Show that $\lim_{r \rightarrow 0} G_s(\mathbf{r}) = \frac{n}{2}$ and $\lim_{r \rightarrow \infty} G_s(\mathbf{r}) = 0$.

Likewise, one can define the *pair correlation function* $g_{ss'}(\mathbf{x} - \mathbf{y})$ by

$$\left(\frac{n}{2}\right)^2 g_{ss'}(\mathbf{x} - \mathbf{y}) = \langle \Phi_0 | \Psi_s^\dagger(\mathbf{x}) \Psi_{s'}^\dagger(\mathbf{y}) \Psi_{s'}(\mathbf{y}) \Psi_s(\mathbf{x}) | \Phi_0 \rangle . \quad (6)$$

(b) Rewrite Eq. (6) in the form

$$\left(\frac{n}{2}\right)^2 g_{ss'}(\mathbf{x} - \mathbf{y}) = \frac{1}{V^2} \sum_{\mathbf{k}_1 \mathbf{k}_2 \mathbf{q}_1 \mathbf{q}_2} e^{-i(\mathbf{k}_1 - \mathbf{k}_2) \cdot \mathbf{x}} e^{-i(\mathbf{q}_1 - \mathbf{q}_2) \cdot \mathbf{y}} \langle \Phi_0 | a_{\mathbf{k}_1, s}^\dagger a_{\mathbf{q}_1, s'}^\dagger a_{\mathbf{q}_2, s'} a_{\mathbf{k}_2, s} | \Phi_0 \rangle . \quad (7)$$

(c) Assume first that $s \neq s'$. Calculate $g_{ss'}(\mathbf{x} - \mathbf{y})$.

(d) Now consider the case where $s = s'$ and calculate $g_{ss}(\mathbf{x} - \mathbf{y})$. Plot the quantity $g_{ss}(\mathbf{x} - \mathbf{y})$ as a function of $|\mathbf{x} - \mathbf{y}|$.