

Computational Quantum Physics Exercise 4

Problem 4.1 Tight binding model

In this exercise you will solve the tight-binding model

$$H = \sum_{\langle i,j \rangle, \sigma} \left(t_{ij} c_{i,\sigma}^\dagger c_{j,\sigma} + h.c. \right) \quad (1)$$

on the simple square lattice analytically. You can assume uniform hopping amplitudes $t_{ij} = t$. Assuming the lattice has $N = L \times L$ sites with spacing a and periodic boundary conditions, you replace the field operators by their Fourier transform

$$c_{i,\sigma} = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{x}_i} c_{\mathbf{k},\sigma}, \quad (2)$$

where the lattice momenta run over N points of the Brillouin zone

$$\mathbf{k} = \frac{2\pi}{aL} (n_x, n_y), \quad n_{x,y} = \{0, \dots, L-1\}. \quad (3)$$

Bring the Hamiltonian into the diagonal form

$$H = \sum_{\mathbf{k}, \sigma} \epsilon(\mathbf{k}) n_{\mathbf{k},\sigma} \quad (4)$$

and identify the dispersion relation $\epsilon(\mathbf{k})$.

Problem 4.2 Exact diagonalization of lattice Hamiltonians

The most accurate method for solving a quantum many-body problem is exact diagonalization of the Hamiltonian matrix. In this exercise we will have a close look at the first two steps of this task: For different models on a two-site lattice we will construct the Hamiltonian matrix and break it down into blocks corresponding to different symmetry sectors.

For each of the models described below

- write down the full Hamiltonian matrix H in a sensible basis,
- partition the basis states into sectors such that H matrix elements between states from different sectors vanish,
- diagonalize the block of H belonging to each sector.

What are the physical symmetries that allow you to split the Hamiltonians into blocks?

Models

1. Spin-1/2 Heisenberg model

$$H = J\vec{S}_1 \cdot \vec{S}_2 \quad (5)$$

for $J = \pm 1$. Here, the spin operators can be written in terms of the Pauli matrices $\vec{S} = \frac{\hbar}{2}(\sigma_x, \sigma_y, \sigma_z)^T$.

2. Spin-1 Heisenberg model

Same Hamiltonian as above, but this time the \vec{S}_i are Spin-1 operators. *Hint:* Express $S_{x,y}$ in terms of the ladder operators S_{\pm} as demonstrated in the lecture script.

3. Bose-Hubbard model

$$H = -t \left(b_1^\dagger b_2 + b_2^\dagger b_1 \right) + \frac{U}{2} \sum_{i=1}^2 n_i (n_i - 1) \quad (6)$$

Now the two sites can be occupied by spinless bosons ($b_i^{(\dagger)}$: bosonic annihilation (creation) operators). As each site could hold an arbitrary number of bosons, you have to limit the total number of particles, e.g. to $N_{max} = 4$. Fix $t = 1$ and diagonalize the system for $U = -1, 1, 2$. Which cases does the particle number cut-off seem reasonable for?

4. (Fermi-) Hubbard model

This is the regular Hubbard model with fermions carrying a spin-1/2 as described in in chapter 6 of the lecture script.

5. $t - J$ model

This Hamiltonian is also introduced in the script.