Metals - electron-hole excitations

linear response function:

$$\chi_0(\vec{q},\omega) = \frac{1}{\Omega} \sum_{\vec{k}\,,s} \frac{n_{0\,\vec{k}+\vec{q}\,,s} - n_{0\,\vec{k}\,,s}}{\epsilon_{\,\vec{k}+\vec{q}} - \epsilon_{\,\vec{k}} - \hbar\omega - i\hbar\eta} \qquad \text{Lindhard}$$
function

Lindhard

$$\lim_{\eta \to 0_{+}} \frac{1}{z - i\eta} \\ = \mathcal{P}\left(\frac{1}{z}\right) + i\pi\delta(z) \\ = \begin{cases} \chi_{01}(\vec{q}, \omega) = \frac{1}{\Omega} \sum_{\vec{k}, s} \mathcal{P}\left(\frac{n_{0, \vec{k} + \vec{q}} - n_{0, \vec{k}}}{\epsilon_{\vec{k} + \vec{q}} - \epsilon_{\vec{k}} - \hbar\omega}\right) \\ \chi_{02}(\vec{q}, \omega) = \frac{1}{\Omega} \sum_{\vec{k}, s} (n_{0, \vec{k} + \vec{q}} - n_{0, \vec{k}})\delta(\epsilon_{\vec{k} + \vec{q}} - \epsilon_{\vec{k}} - \hbar\omega) \\ E = \epsilon_{\vec{k} + \vec{q}} - \epsilon_{\vec{k}} \\ k + q, s' \\ \end{cases}$$

linear response function:

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 function

Lindhard

$$\begin{split} \text{small-}q\text{-limit:} \qquad & \epsilon_{\,\vec{k}\,+\,\vec{q}} \approx \epsilon_{\,\vec{k}}\,+\,\vec{q}\,\cdot\,\vec{\nabla}_{\,\vec{k}}\,\,\epsilon_{\,\vec{k}} = \epsilon_{\,\vec{k}}\,+\,\vec{q}\,\cdot\,\hbar\,\vec{v}_{\,\,\vec{k}} \\ n_{0,\,\vec{k}\,+\,\vec{q}} \approx & n_{0,\,\vec{k}}\,+\,\frac{\partial n_0}{\partial \epsilon}\,\vec{q}\,\cdot\,\vec{\nabla}_{\,\vec{k}}\,\,\epsilon_{\,\vec{k}} = n_{0,\,\vec{k}}\,-\,\delta(\epsilon_{\,\vec{k}}\,-\epsilon_F)\,\vec{q}\,\cdot\hbar\,\vec{v}_{\,\,\vec{k}} \end{split}$$

$$\chi_0(\vec{q},\omega) \approx -2 \int \frac{d^3k}{(2\pi)^3} \frac{\vec{q} \cdot \vec{v}_F \delta(\epsilon_{\vec{k}} - \epsilon_F)}{\vec{q} \cdot \vec{v}_F - \omega - i\eta}$$
$$\approx \frac{n_0 q^2}{m(\omega + i\eta)^2} \left(1 + \frac{3}{5} \frac{v_F^2 q^2}{(\omega + i\eta)^2}\right)$$

$$\chi_0(\vec{q},\omega) \approx \frac{n_0 q^2}{m(\omega + i\eta)^2} \left(1 + \frac{3}{5} \frac{v_F^2 q^2}{(\omega + i\eta)^2} \right) = \frac{n_0 q^2}{m(\omega + i\eta)^2} R(q,\omega)^2$$

$$\varepsilon(\vec{q},\omega) = 1 - \frac{4\pi e^2}{q^2} \chi_0(\vec{q},\omega) \approx 1 - \frac{\omega_p^2}{\omega^2}$$
 plasma frequence
$$\omega_p^2 = \frac{4\pi e^2 n_0}{m}$$

renormalized: $\chi(\vec{q},\omega) = \frac{\chi_0(\vec{q},\omega)}{1 - \frac{4\pi e^2}{e^2}\chi_0(\vec{q},\omega)}$

$$1.2 - 4\pi e^{2}$$

plasma frequency

$$\begin{split} \chi(\vec{q},\omega) &\approx \frac{n_0 q^2 R(q,\omega)^2}{(\omega + i\eta)^2 - \frac{4\pi e^2 n_0^2}{m} R(q,\omega)^2} \\ &= \frac{n_0 q^2 R(q,\omega)}{\omega_p} \left\{ \frac{1}{\omega + i\eta - \omega_p R(q,\omega)} - \frac{1}{\omega + i\eta + \omega_p R(q,\omega)} \right\} \end{split}$$

$$\chi(\vec{q},\omega) \approx \frac{n_0 q^2 R(q,\omega)}{\omega_p} \left\{ \frac{1}{\omega + i\eta - \omega_p R(q,\omega)} - \frac{1}{\omega + i\eta + \omega_p R(q,\omega)} \right\}$$

spectrum in imagnary part

$$\chi''(\vec{q},\omega) \approx \frac{\pi n_0 q^2 R(q,\omega_p)}{\omega_p} \left[\delta(\omega - \omega_p R(q,\omega_p)) - \delta(\omega + \omega_p R(q,\omega_p)) \right]$$

$$\omega(\vec{q}) = \omega_p R(q, \omega_p)$$

$$= \omega_p \left\{ 1 + \frac{3v_F^2 q^2}{10\omega_p^2} + \cdots \right\}$$

 $\frac{\hbar\omega}{4\epsilon_F}$ $\frac{\hbar\omega}{1}$ plasmon
electron-hole-excitations
finite etime $\frac{1}{a/2k_F}$

Landau damping in e-h-continuum decay of plasmon into electron-hole excitations finite lifetime

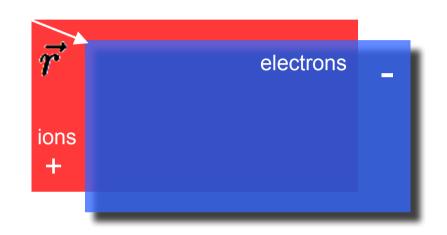
$$\omega_p^2 = \frac{4\pi e^2 n_0}{m}$$

simple metals
$$n_0=rac{3}{4\pi d^3}=rac{3}{4\pi (r_s a_B)^3}$$

$$r_{s,Li} = 3.22$$
 $r_{s,Na} = 3.96$ $r_{s,K} = 4.86$

metal	$\omega_p^{(\mathrm{exp})} \; \; [\mathrm{eV}]$	$\omega_p^{(\mathrm{theo})} \; \; [\mathrm{eV}]$
Li	7.1	8.5
Na	5.7	6.2
K	3.7	4.6
Mg	10.6	_
Al	15.3	-

classical picture



restoring force

$$ec{P} = -e n_0 \, ec{r}$$
 $ec{E} = -4 \pi \, ec{P}$

$$ec{E} = -4\pi\,ec{P}$$

equation of motion

$$m rac{d^2 \, ec{r}}{dt^2} = -e \, ec{E} \, = -4 \pi e^2 n_0 \, ec{r}$$

harmonic oscillator

oscillation frequency

$$\omega_p^2 = rac{4\pi e^2 n_0}{m}$$

plasma frequency